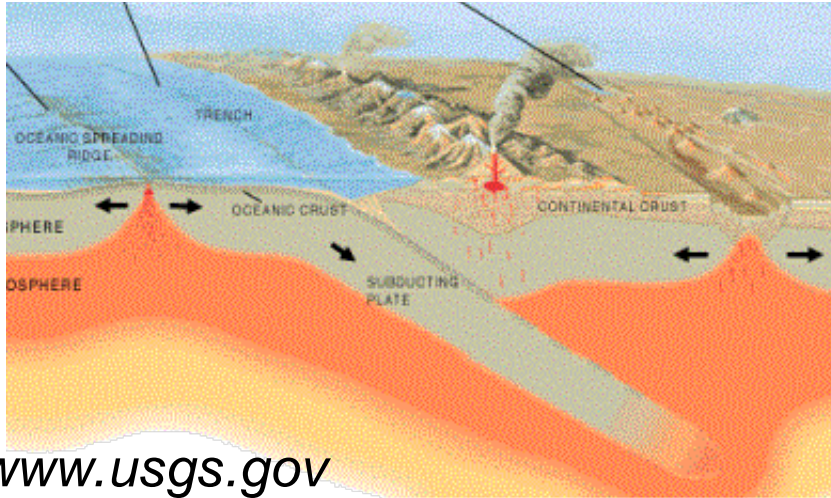


Seismic tomography: Art or science?

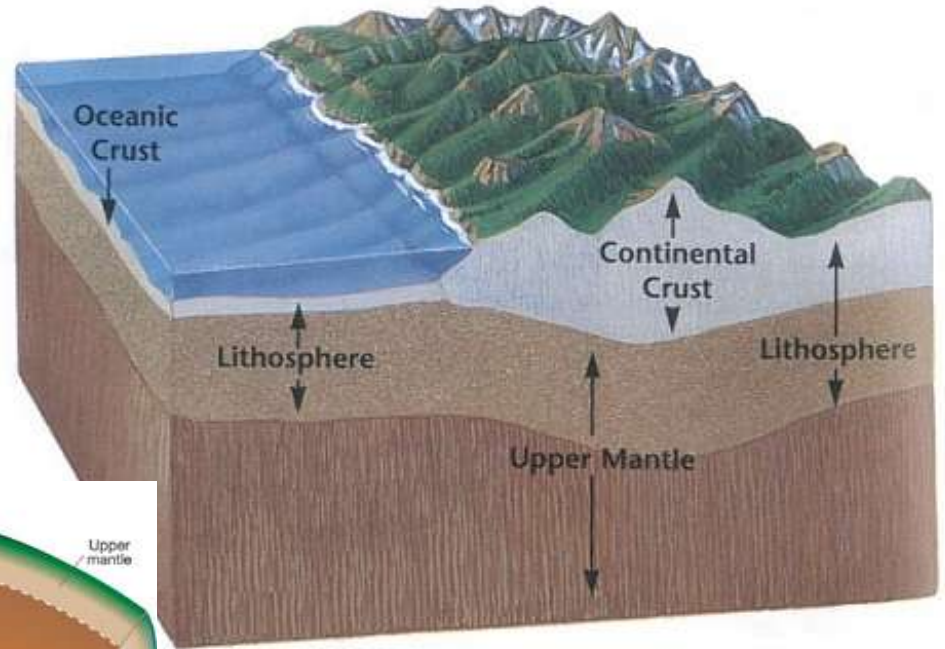
Frederik J Simons
Princeton University



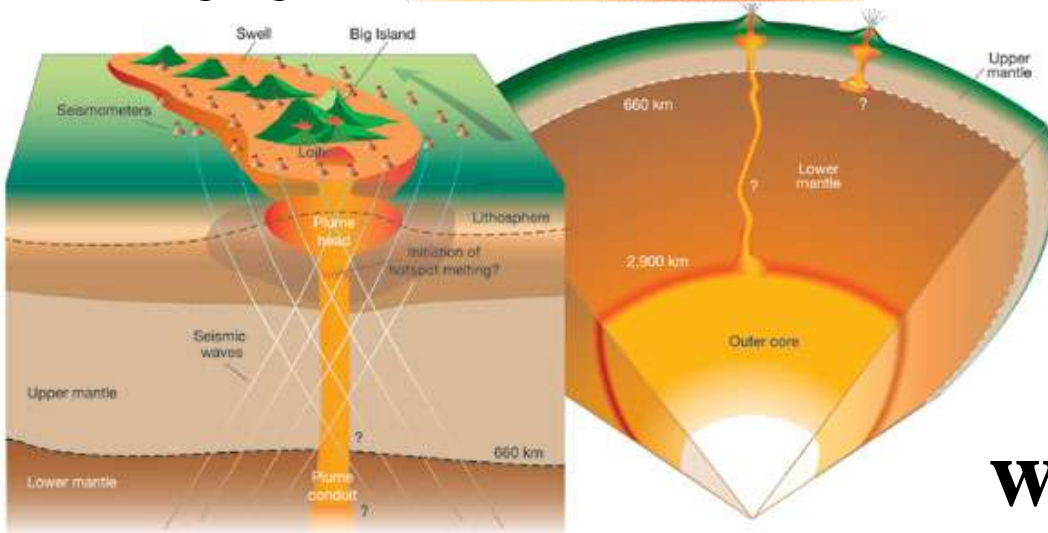
What's inside the Earth?



www.usgs.gov



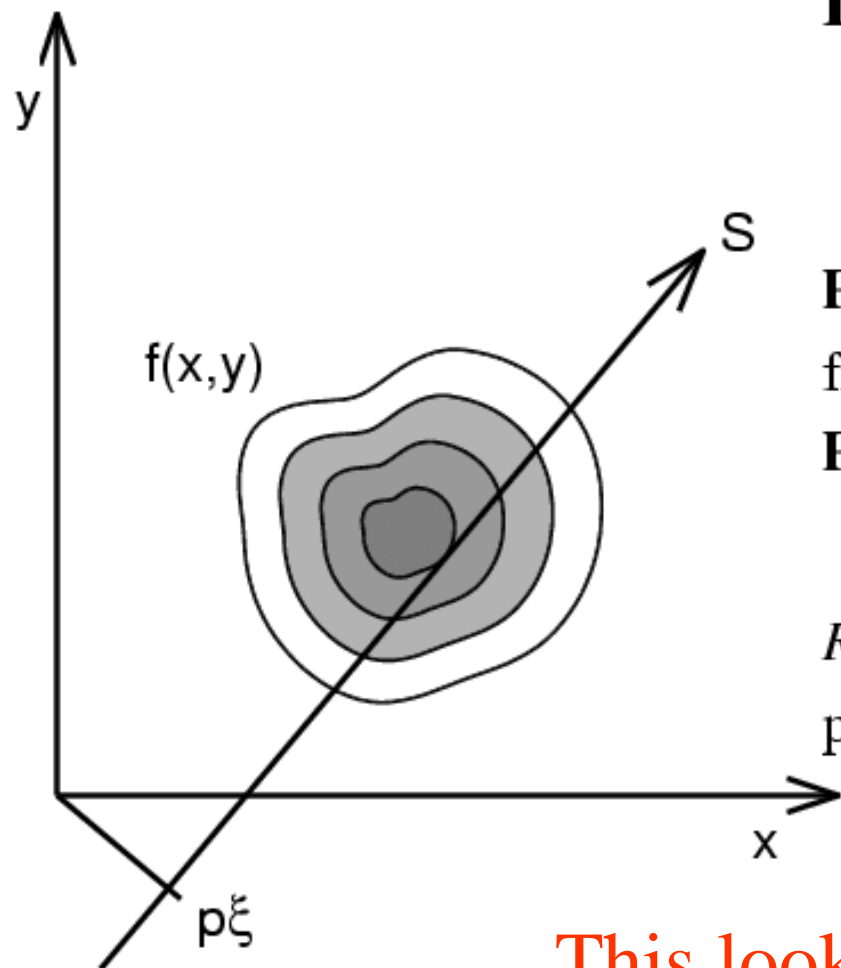
www.glencoe.com



Dalton, *Nature*
2003

Only seismic waves have actually been there, done that

The seismic tomography problem



Inverting the Radon transform

$$\mathcal{R}[f(x, y)](p, \xi) = \int_S f(x, y) ds \quad (1)$$

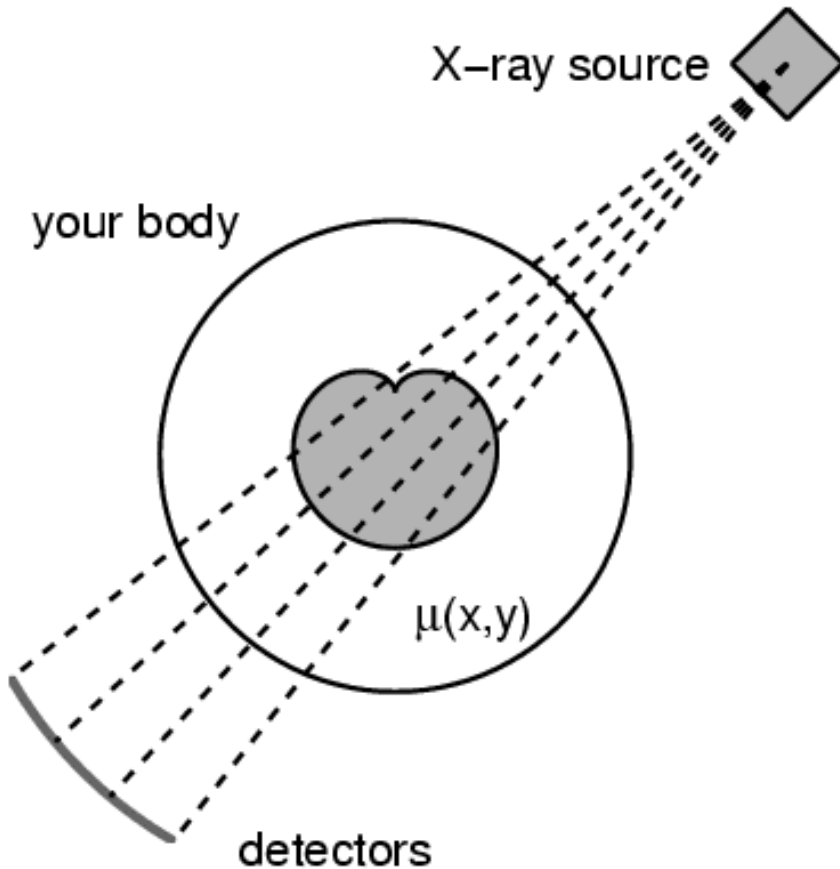
Purpose: Reconstruction of functions from their line integrals (projections).

Problem: Given $\mathcal{R}[f(x, y)](p, \xi)$, find $f(x, y)$.

Radon [1917] derived a solution to this problem, giving an expression for \mathcal{R}^{-1} .

This looks more complicated than it is; and that's my point.

What is $f(x, y)$? Medical applications.



X-ray absorption & scattering

Tissues and bones have \neq absorption and scattering coefficients $\mu(x, y)$.

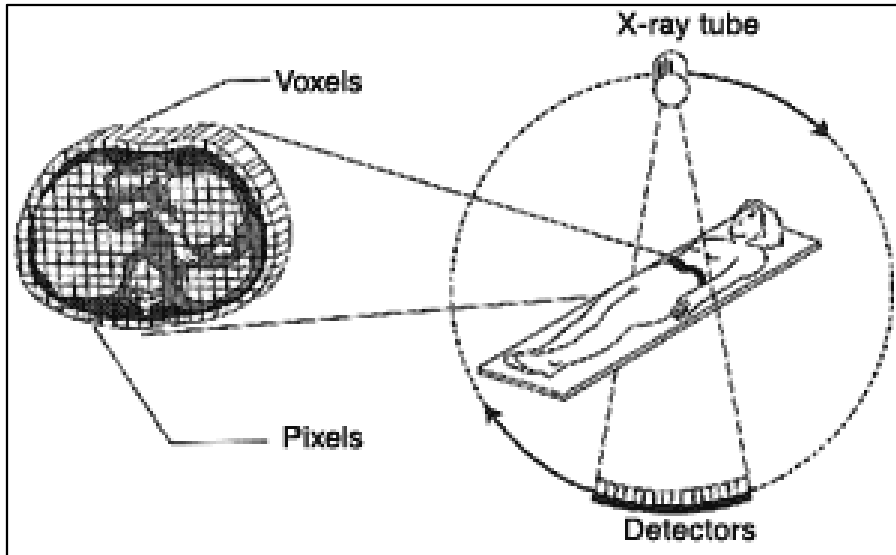
Recorded intensity goes as

$$I = I_0 \exp \left[\int_{\text{ray}} -\mu(x, y) ds \right]. \quad (2)$$

Sources and detectors rotate to achieve perfect "coverage".

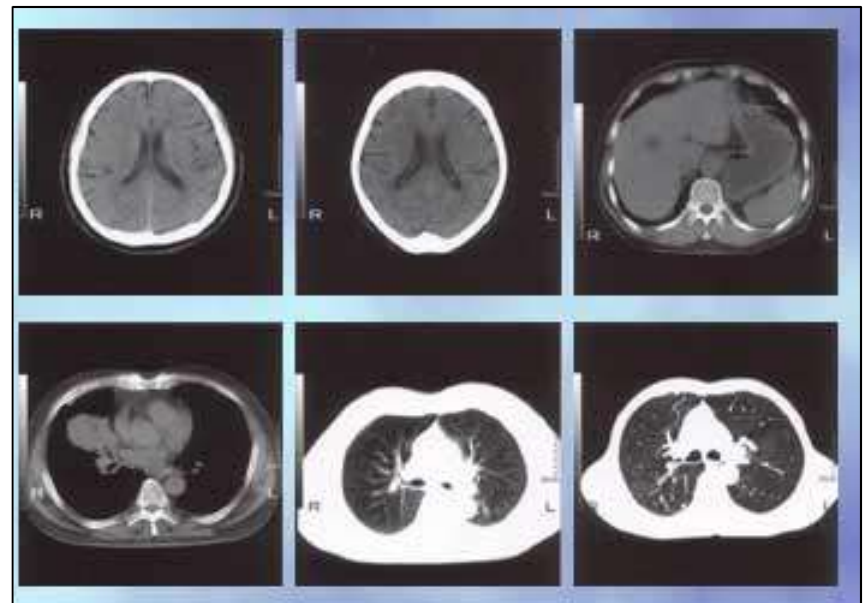
This looks simpler than it is;
and that's my point.

X-Ray attenuation tomography

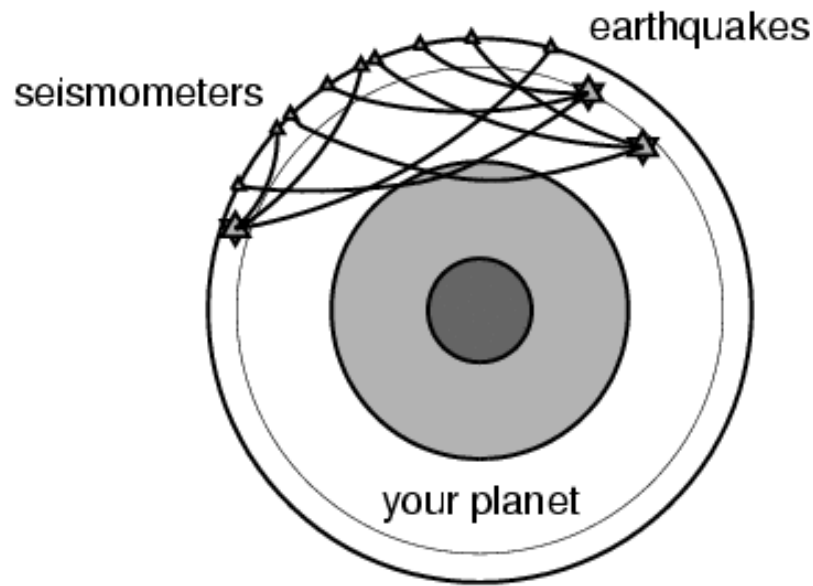


Projections from all angles:
X-ray intensity

Reconstructed image:
X-ray attenuation constants



What is $f(x, y, z)$? Seismic wavespeeds.



Travel-time tomography

The Earth is made of a heterogeneity of seismic velocities $v(x, y, z)$.

Travel-time anomalies go as

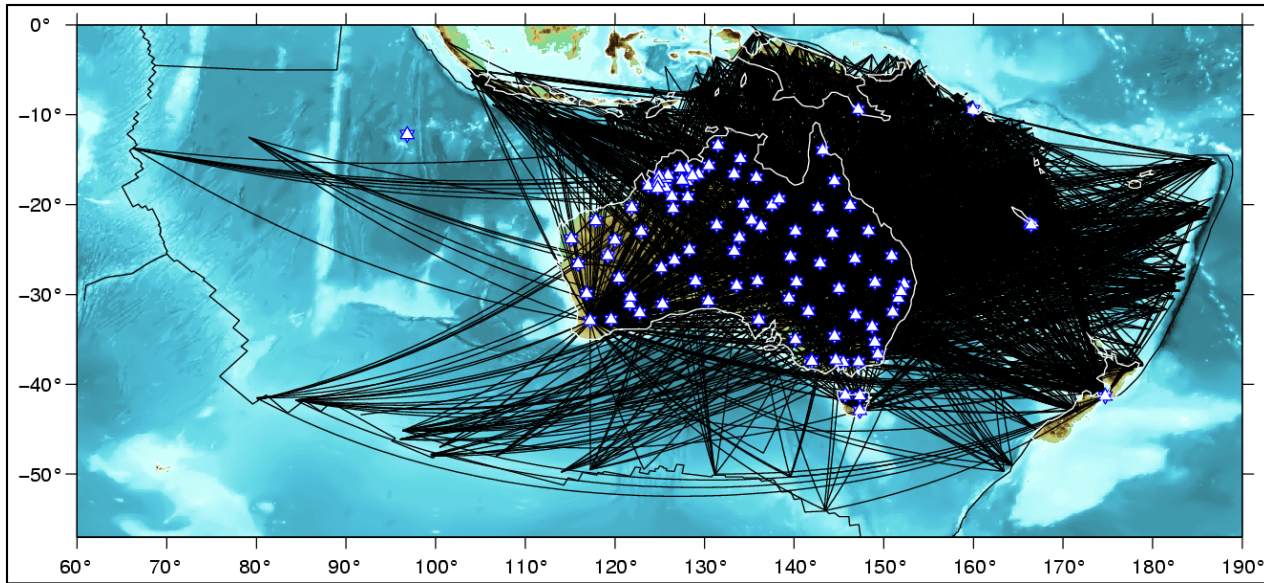
$$\delta t = \int_{\text{ray}} \frac{1}{\delta v(x, y, z)} ds. \quad (3)$$

Waveform tomography

Arrival times depend on the wavelength of the seismic phases.

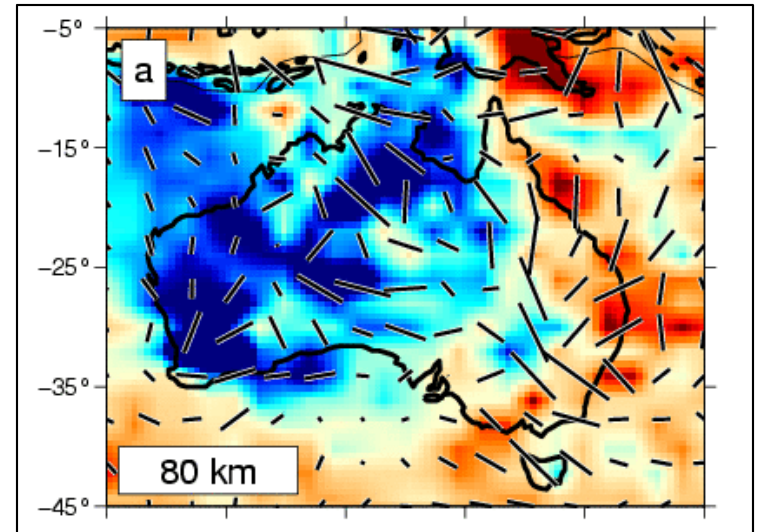
All raypaths curve and coverage is far from perfect.

Seismic wavespeed tomography



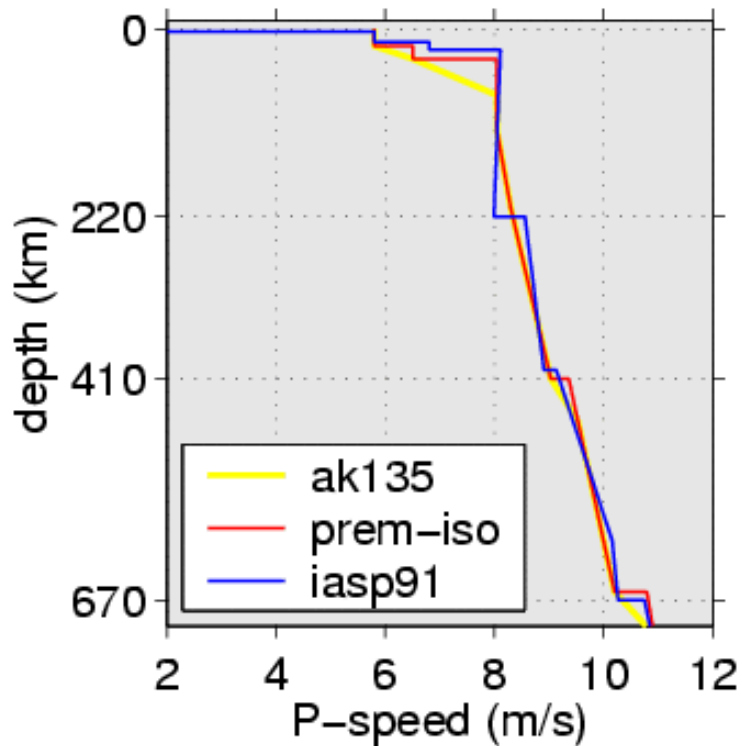
Projections from all angles:
Waveforms and arrival times

Reconstructed image:
Wavespeed variations

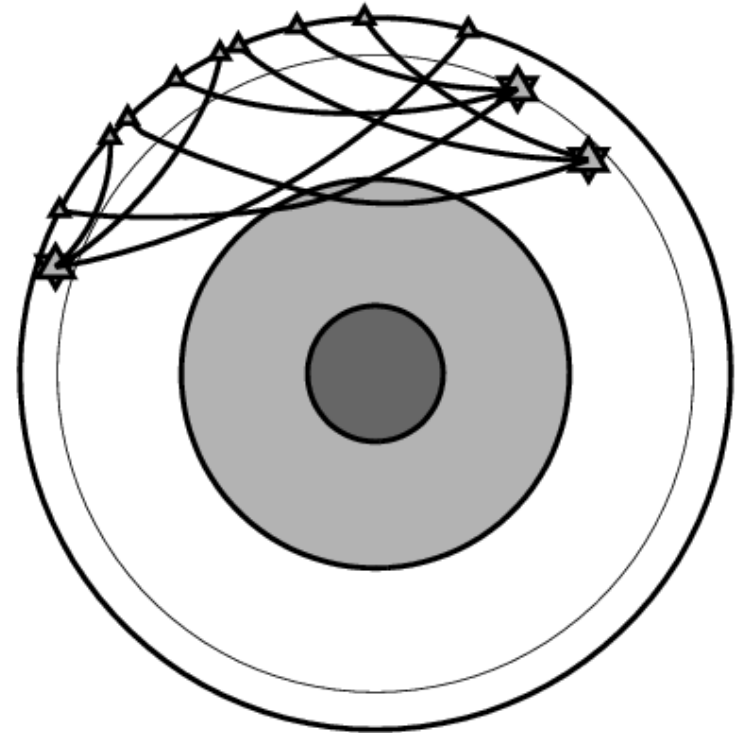


Forward modeling of the wave field, Part I:

Ray tracing, **most 1-D**



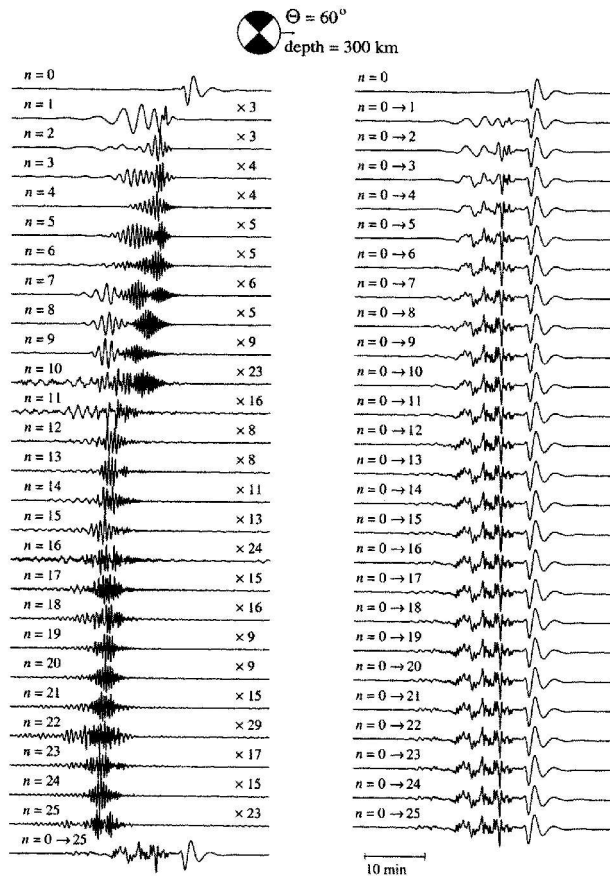
Before



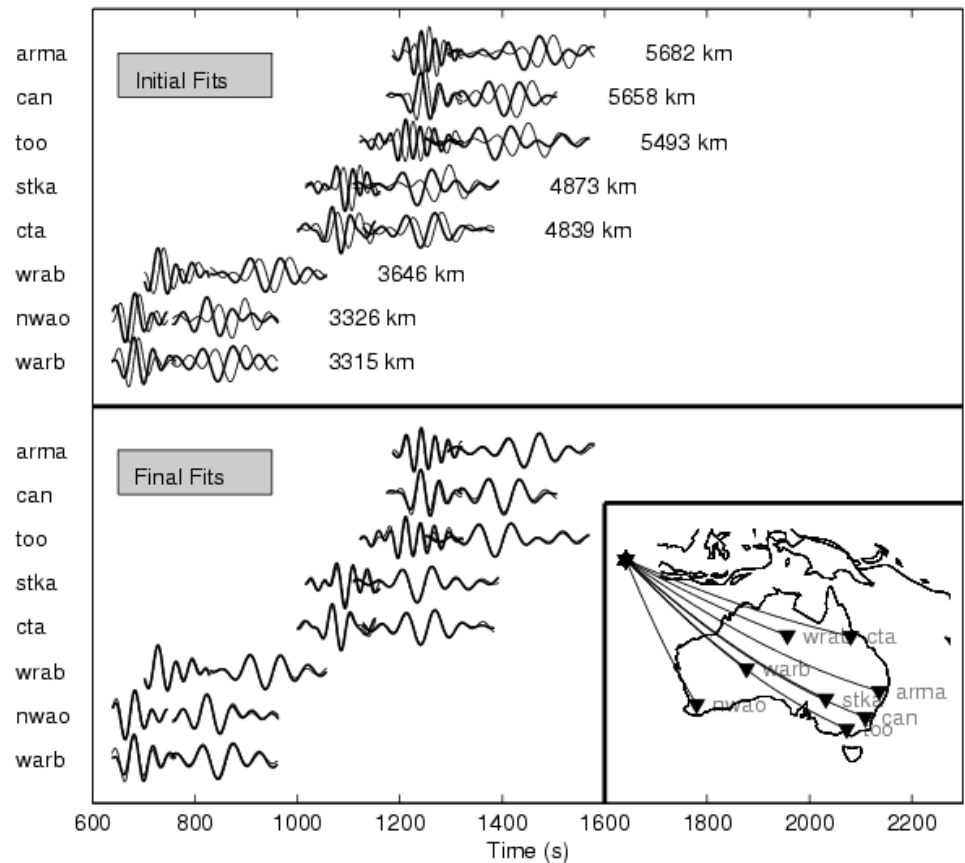
After

Forward modeling of the wave field, Part II:

Normal-mode summation, 1-D



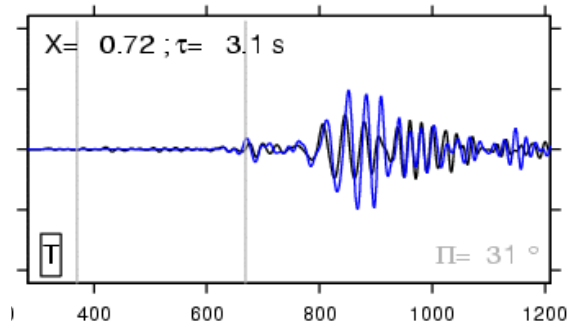
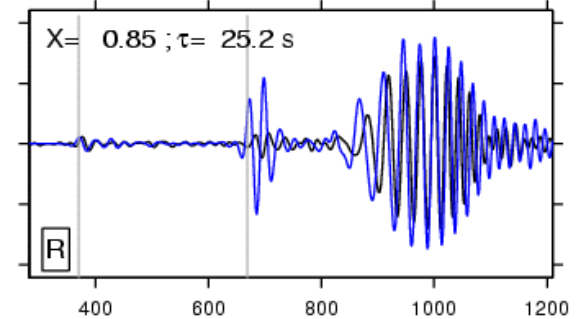
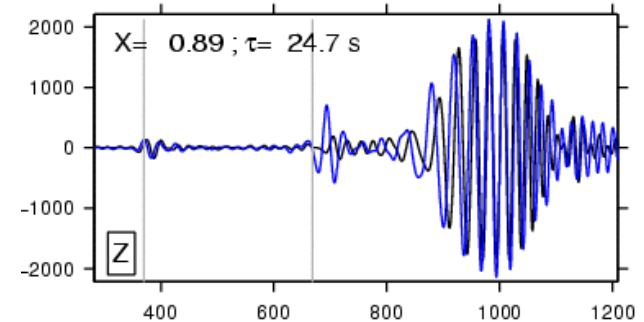
Before



After

Forward modeling of the wave field, Part III: Spectral-element methods, 3-D

Before



STKA, Δ = 30.46°

After

That's all there is to it. Gooby!

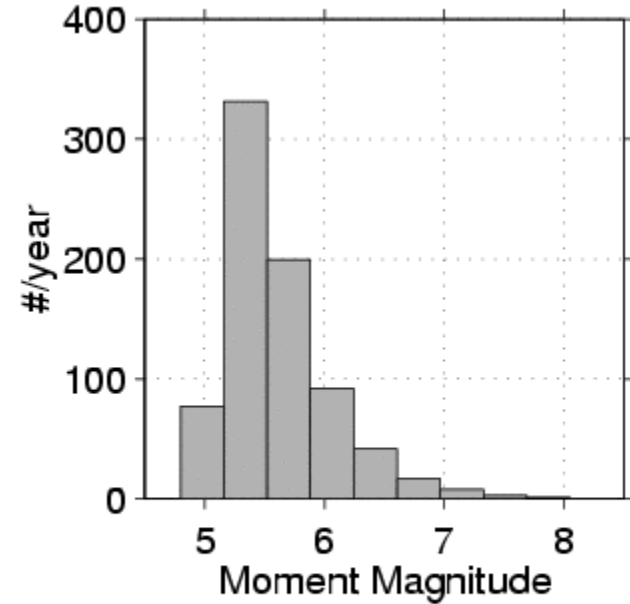
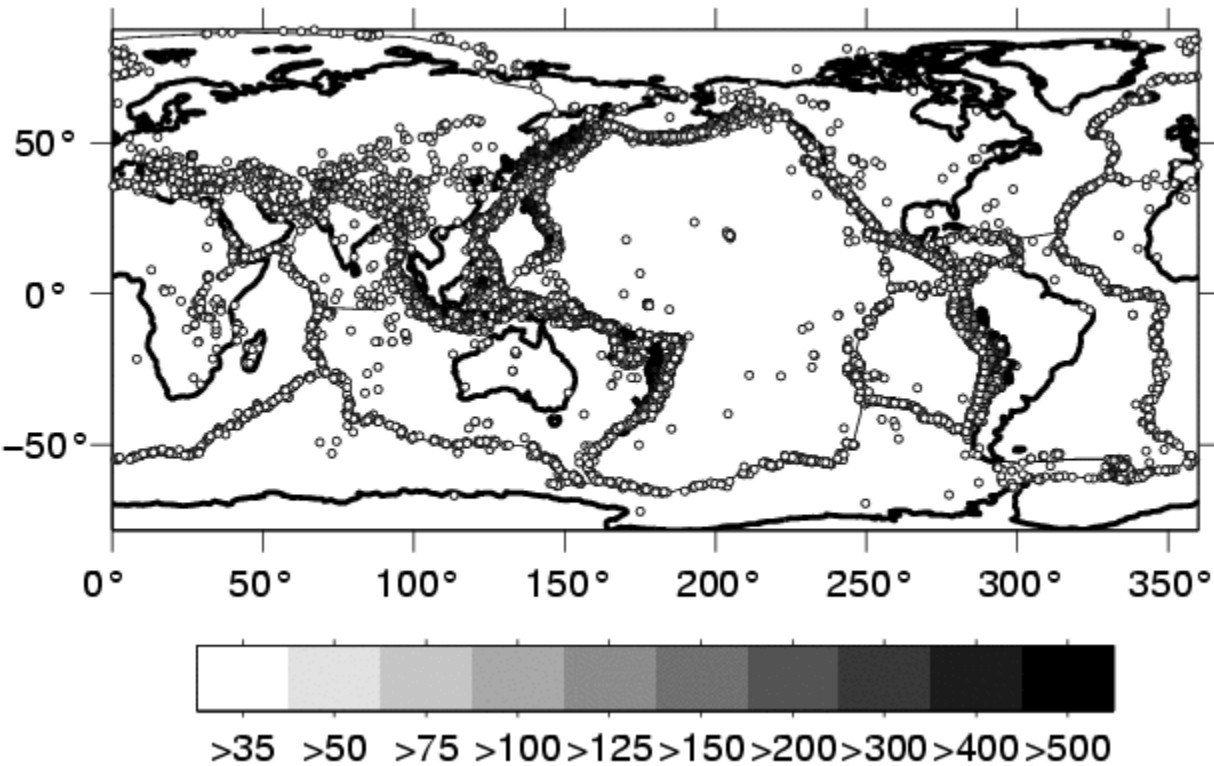
Except:

- **X-ray**: exponential of a line integral
 - **S-ray**: raypath itself is a function of velocity
- } **non-linear functions!**
- Earth **coverage** is non-continuous
 - “Experiment” is done by nature and **not repeatable**
 - Earthquake **source parameters** (location, time) is uncertain

Remedy:

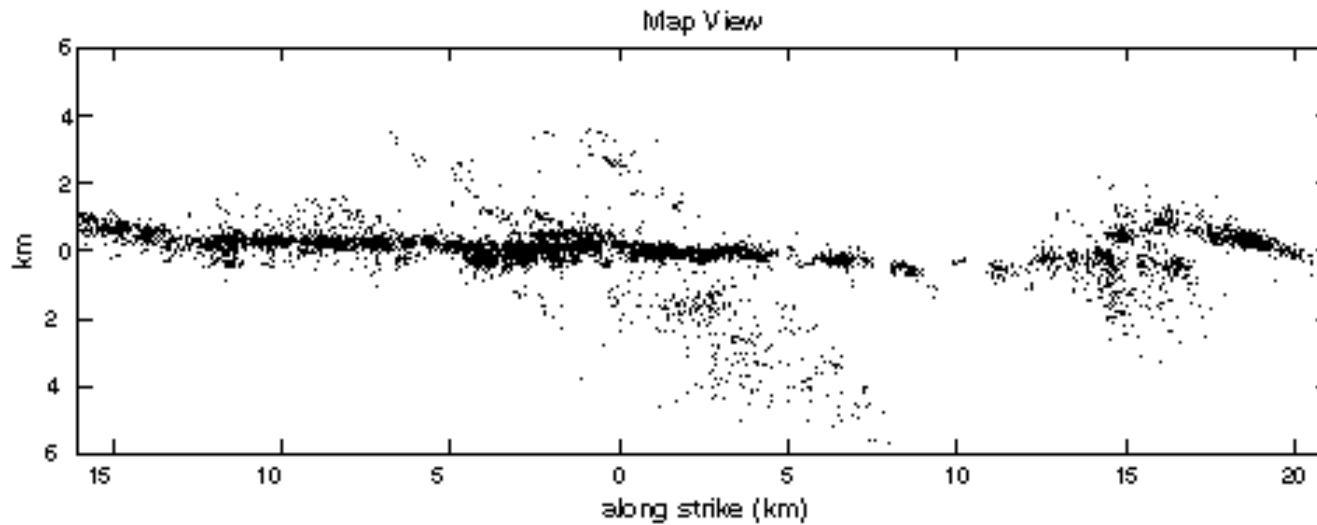
- Linearization
- Discretization
- Regularization (*a priori* information)

Non-continuous **source coverage**

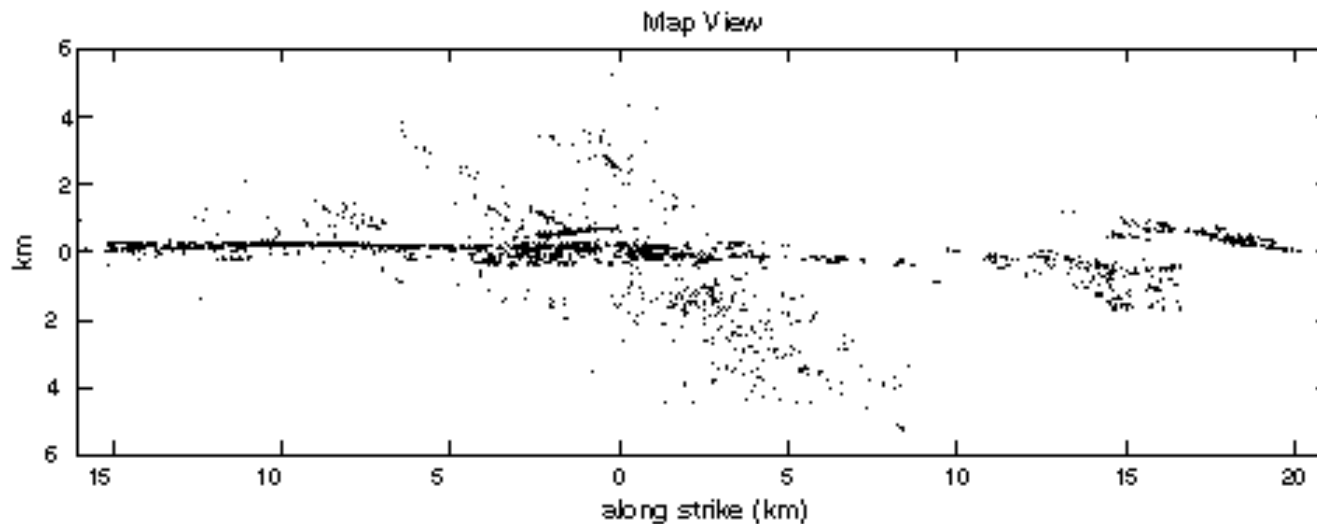


The CMT catalog of large events

Source location – (in)extricably linked



Before



After

Source relocation is big business.

Recipe, Step 1: Linearize!

X-ray

Approximate $\exp(-x) \approx 1 - x$.

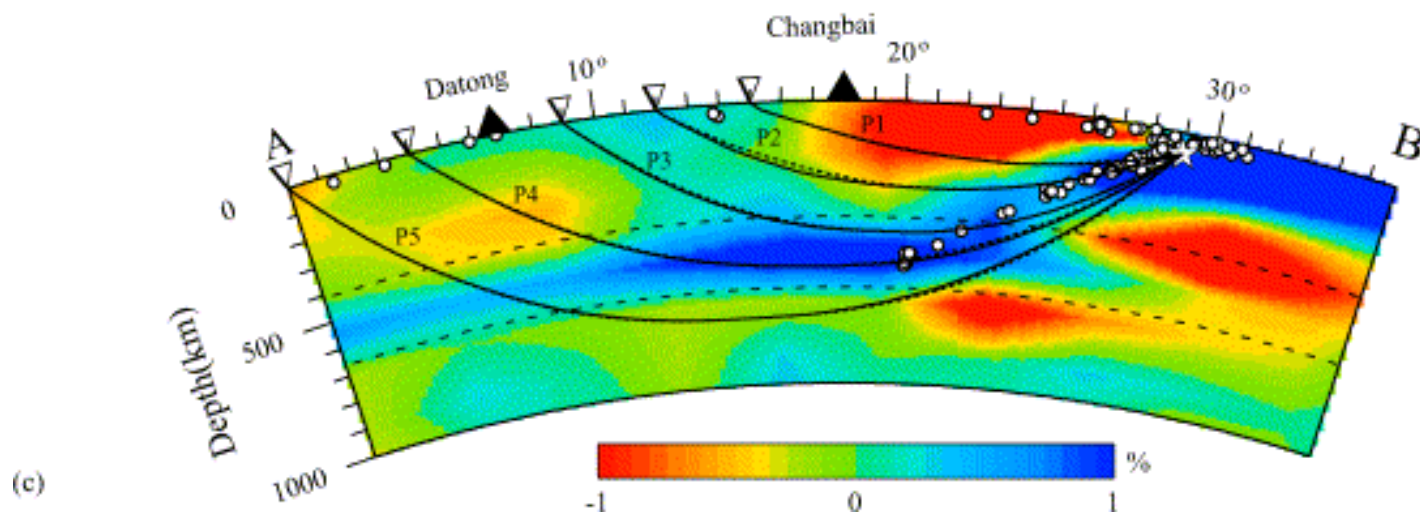
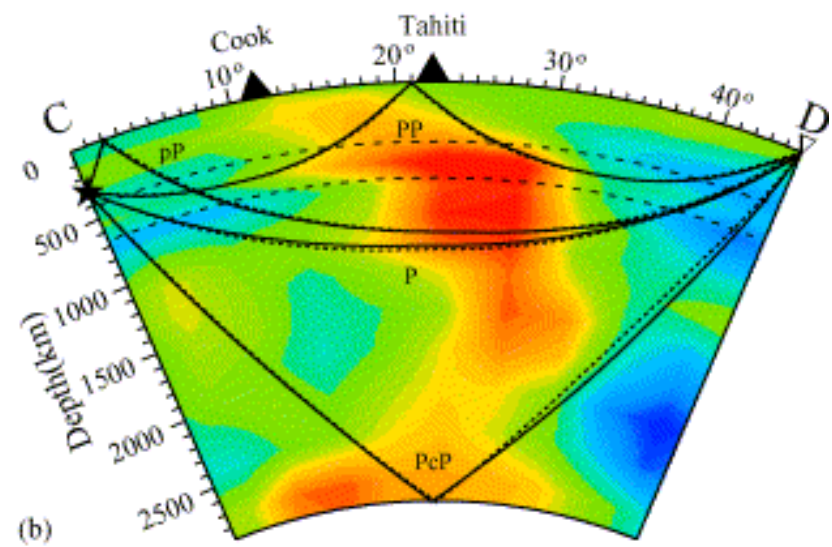
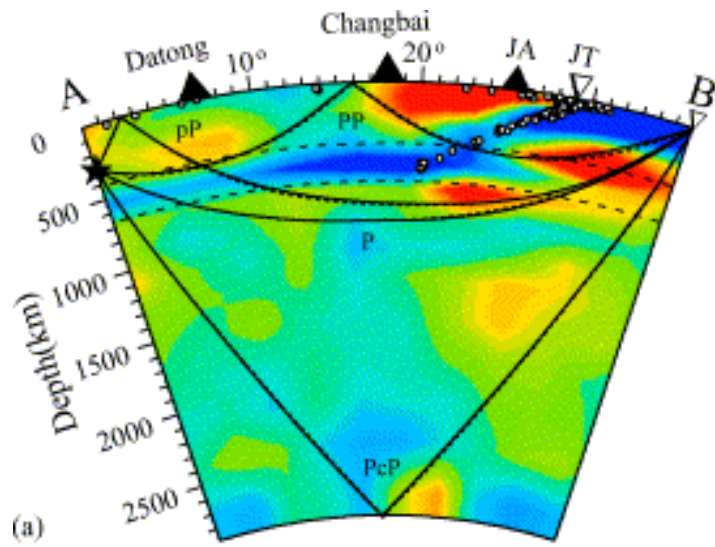
S-ray

Fermat's principle: For a small perturbation of the path, the travel-time (anomaly) is stationary. Using the *slowness*:

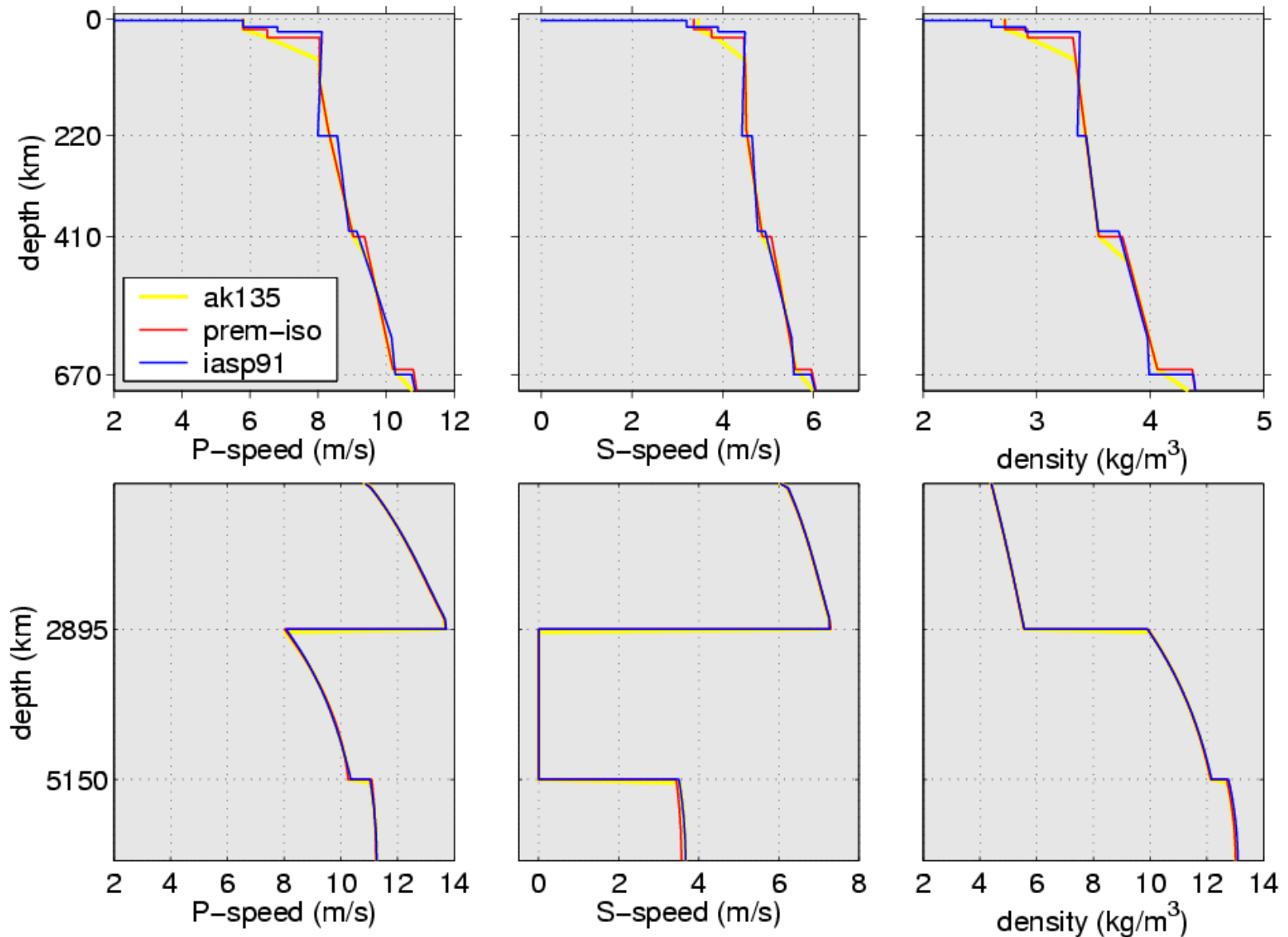
$$\delta s = \frac{1}{\delta v} \rightarrow \delta(\delta t) + \mathcal{O}[(\delta t)^2]. \quad (4)$$

This highlights the importance of the **reference model**, usually a radial model $v(r)$, such as PREM, AK135, IASP91.

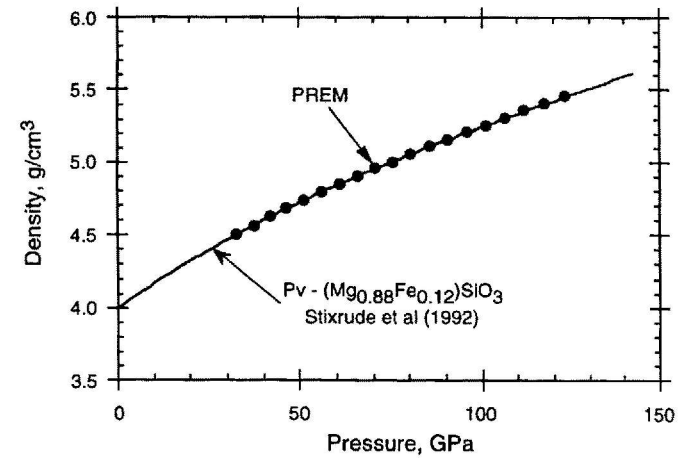
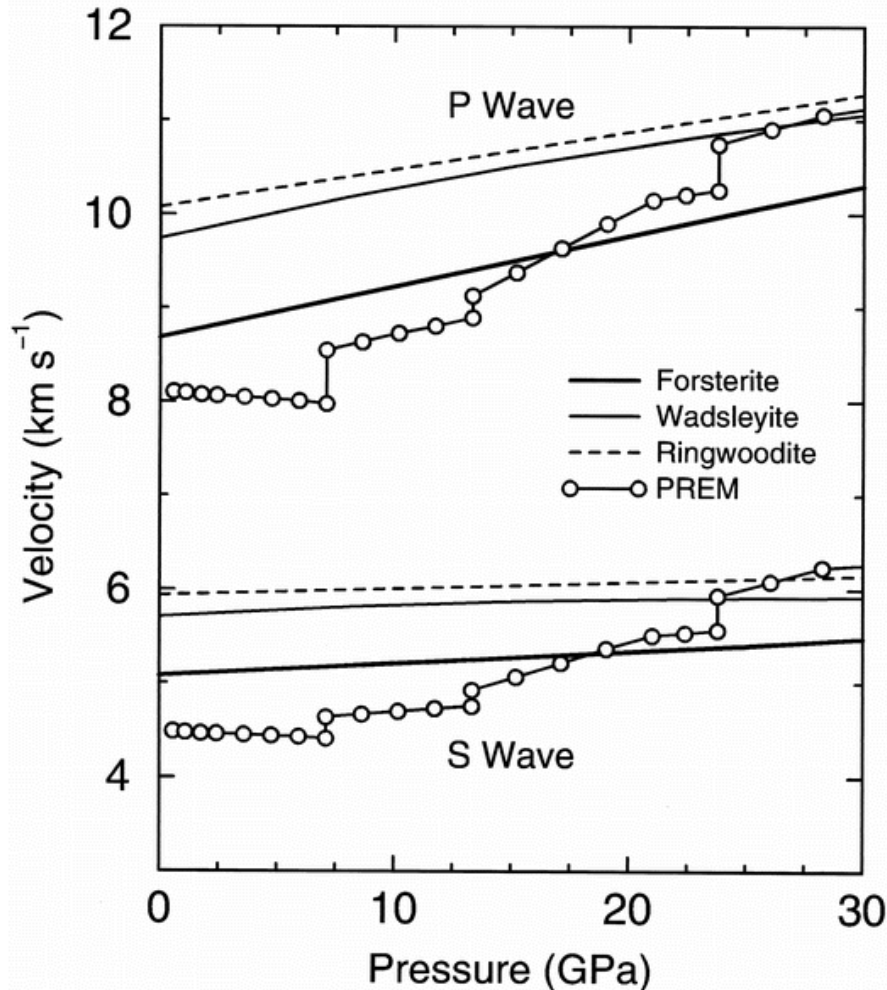
Fermat's Principle at Work for you



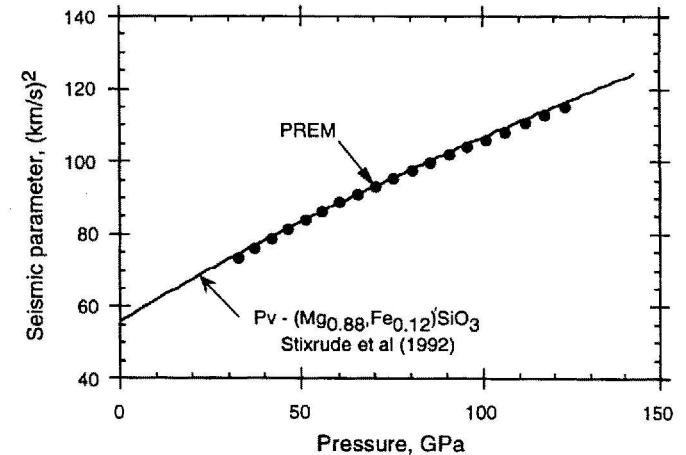
The reference Earth: Radial models



... and at least some of it is true...

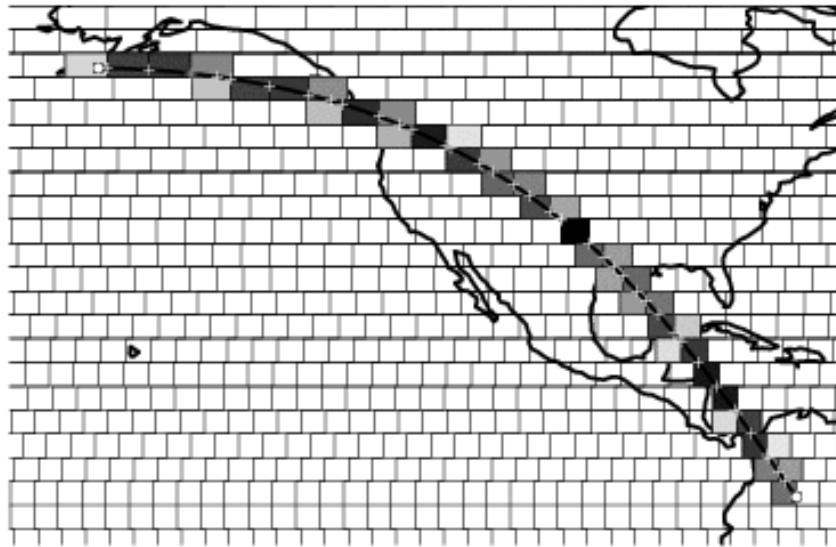


(a)



(b)

Recipe, Step 2: Discretize!



For a set of seismic rays $i = 1 \rightarrow M$, calculate the length spent in each of $j = 1 \rightarrow N$ grid boxes, in each of which it accumulates a proportional fraction of the total travel-time anomaly δt .

$$\delta t_i = L_{ij} \delta s_j \quad \text{or} \quad \delta \mathbf{t} = \mathbf{L} \cdot \delta \mathbf{s} \quad (5)$$

$$\begin{array}{l} \text{M travel-time} \\ \text{anomalies} \end{array} \begin{bmatrix} \vdots \\ \delta t_i \\ \vdots \end{bmatrix} = \begin{array}{c} \begin{bmatrix} \vdots & & \\ \dots & L_{ij} & \dots \\ \vdots & & \end{bmatrix} \\ \text{M} \times \text{N sensitivity matrix} \end{array} \times \begin{array}{c} \begin{bmatrix} \vdots \\ \delta s_j \\ \vdots \end{bmatrix} \\ \text{N slowness} \\ \text{perturbations} \end{array} \quad (6)$$

Letting it simmer: Solving inverse problems

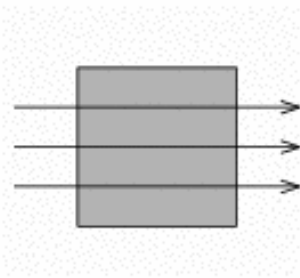
We have: $\mathbf{G} \cdot \mathbf{m} = \mathbf{d}$, which is **linear**.

You think: $\mathbf{m} = \mathbf{G}^{-1} \cdot \mathbf{d}$, but we **can't invert** a non-square $M \times N$ matrix.

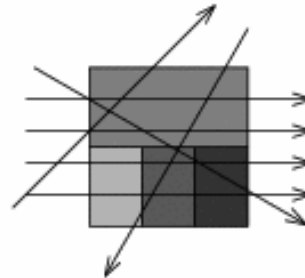
You think: $\mathbf{G}^T \cdot \mathbf{G}$ is square, let's solve $\mathbf{G}^T \cdot \mathbf{G} \cdot \mathbf{m} = \mathbf{G}^T \cdot \mathbf{d}$.

You try: $\mathbf{m} = (\mathbf{G}^T \cdot \mathbf{G})^{-1} \cdot \mathbf{G}^T \cdot \mathbf{d}$.

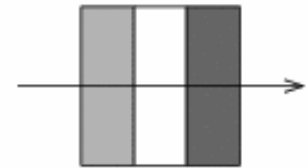
Alas! $\mathbf{G}^T \cdot \mathbf{G}$ may be singular, ill-conditioned, under/over-determined, have (near-)zero eigenvalues, and thus be **not-invertible**.



over-determined, $M > N$



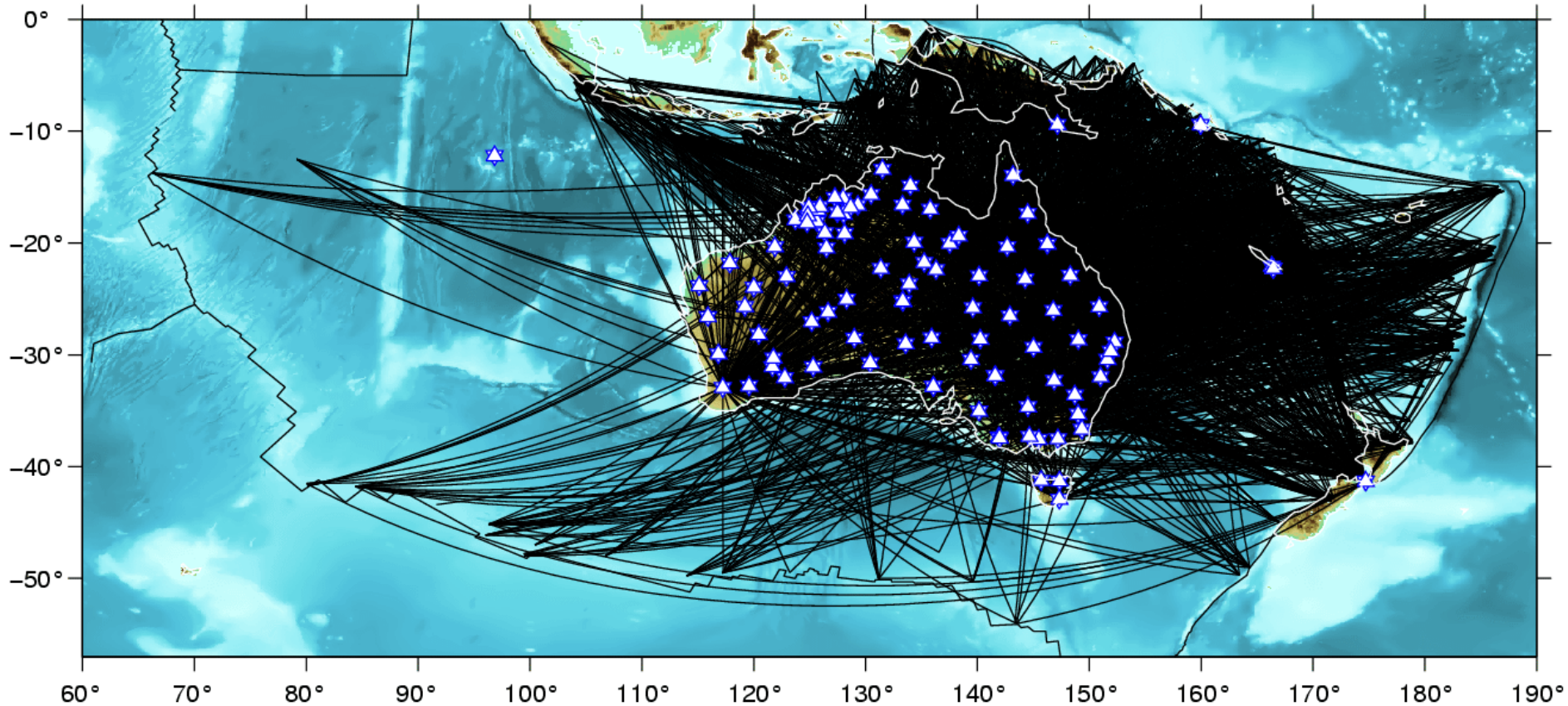
mixed-determined



under-determined, $M < N$

Receiver coverage

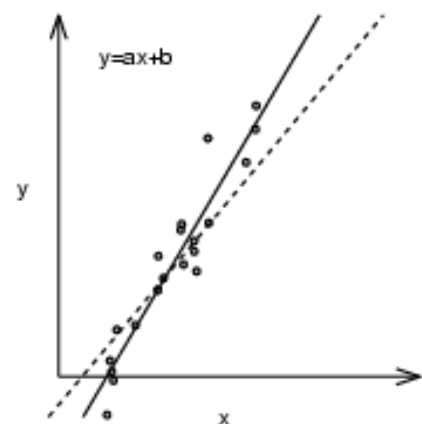
Picking the right continent



A dense path coverage minimizes the amount of a priori information needed

Recipe, Step 3: Regularize!

Over-determined: *More data than unknowns*



Define a *penalty function* Φ on the *error* \mathbf{e} ,
and minimize, by least-squares:

$$\Phi = [\mathbf{G} \cdot \mathbf{m} - \mathbf{d}]^2 = \mathbf{e}^T \cdot \mathbf{e} \quad \text{by} \quad \frac{\partial \Phi}{\partial m_i} = 0. \quad (7)$$

This is a minimization in the *data space*.

Under-determined: *More unknowns than data*

Add equations that minimize some norm in the *model space*:

$$\Phi = \mathbf{e}^T \cdot \mathbf{e} + \mathbf{m}^T \cdot (\mathbf{A}^T \cdot \mathbf{A}) \cdot \mathbf{m}. \quad (8)$$

If $\mathbf{A} = \mathbf{I}$ the identity matrix \rightarrow minimum model norm: **norm damping**.

If $\mathbf{A} = \mathbf{D}$ a difference matrix \rightarrow minimum-roughness: **smoothing**.

Regularization: the Mathematics

Numerical Methods for the Solution of Ill-Posed Problems

by

A. N. Tikhonov
A. V. Goncharenko
V. V. Stepanov
A. G. Yagola

Moscow State University,
Moscow, Russia



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ILL-POSED AND INVERSE PROBLEMS

DEDICATED TO ACADEMICIAN MIKHAIL MISHKALOVICH LITVINSEV
ON THE OCCASION OF HIS 70th BIRTHDAY



Editors:

V.G. Romanov, S.I. Kabanikhin,
Yu.E. Anikonov and A.L. Bukhgeim

/// VSP ///

Utrecht • Boston

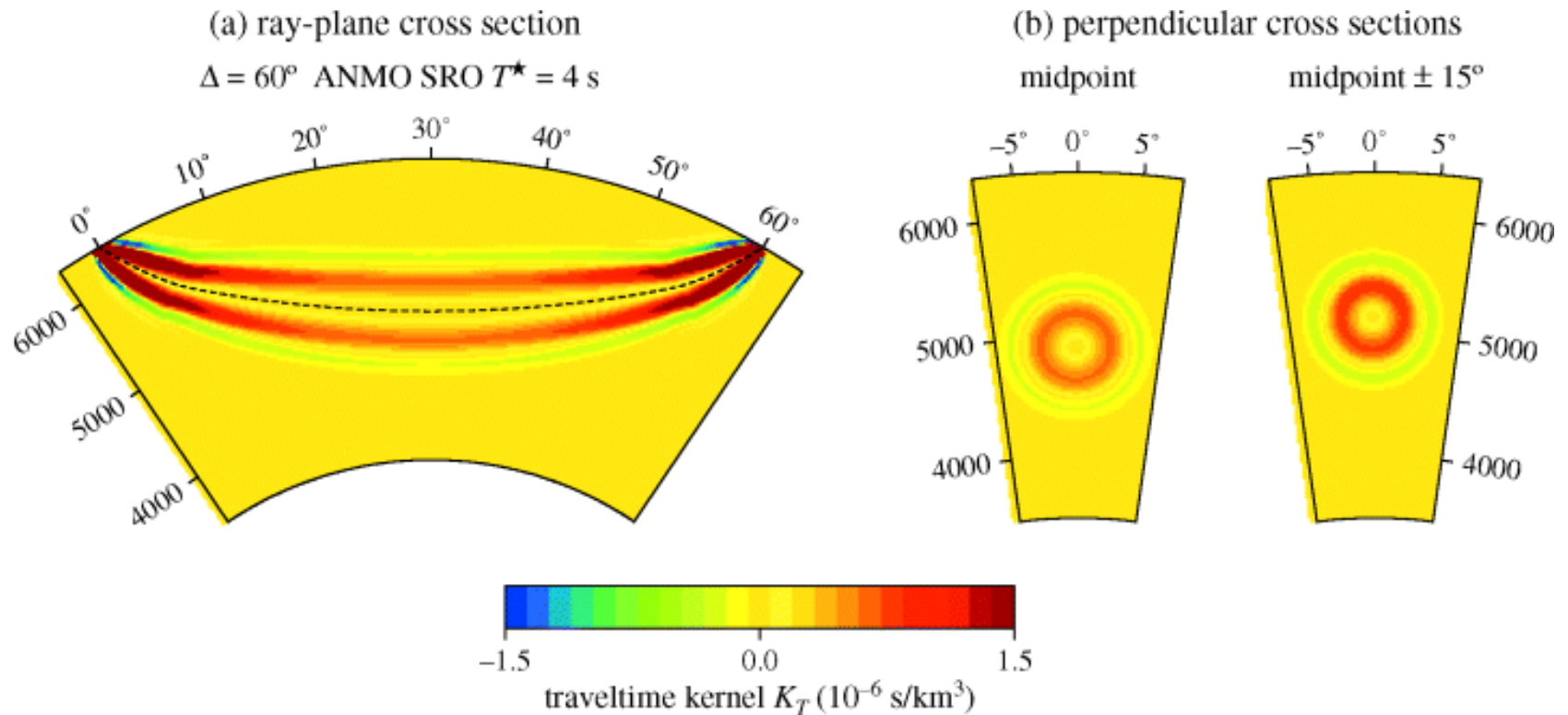
Methods in Geochemistry and Geophysics, 36

GEOPHYSICAL INVERSE THEORY AND REGULARIZATION PROBLEMS

M.S. ZHDANOV

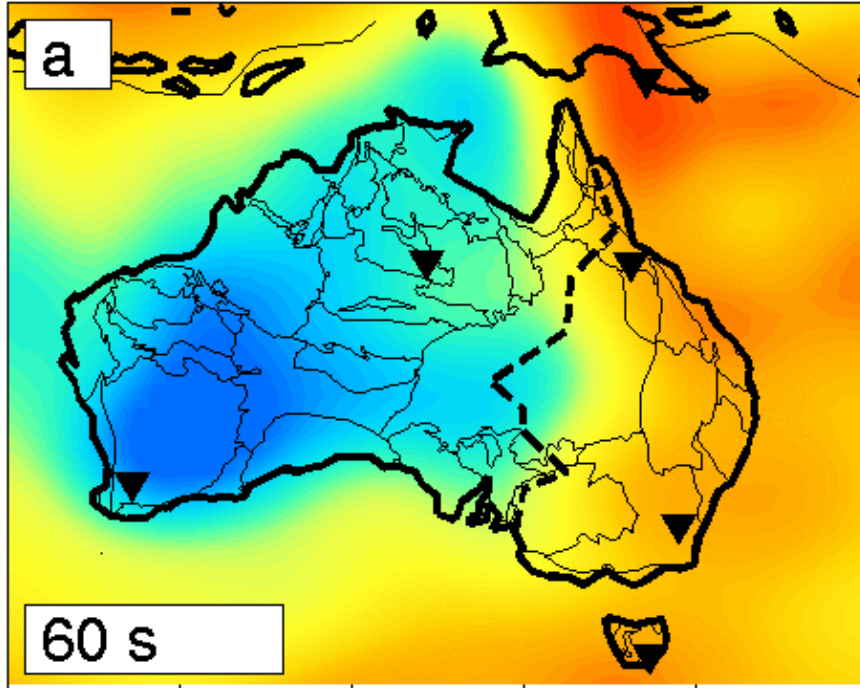
ELSEVIER

Regularization: the Physics

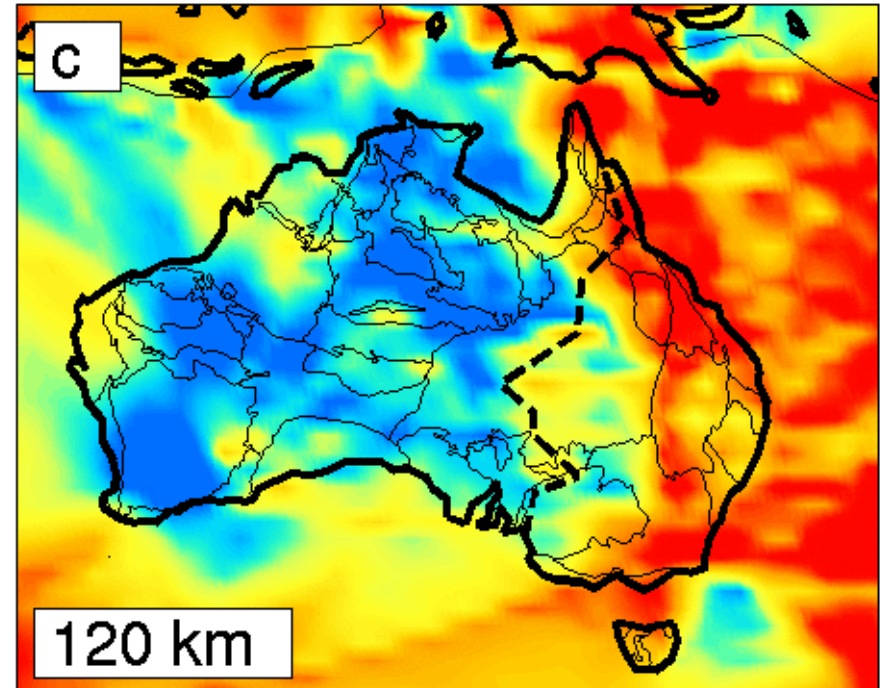


Such “fat” rays sample more of the Earth and thus we need fewer of them to have a well-constrained tomographic problem.

Regularization: the Art

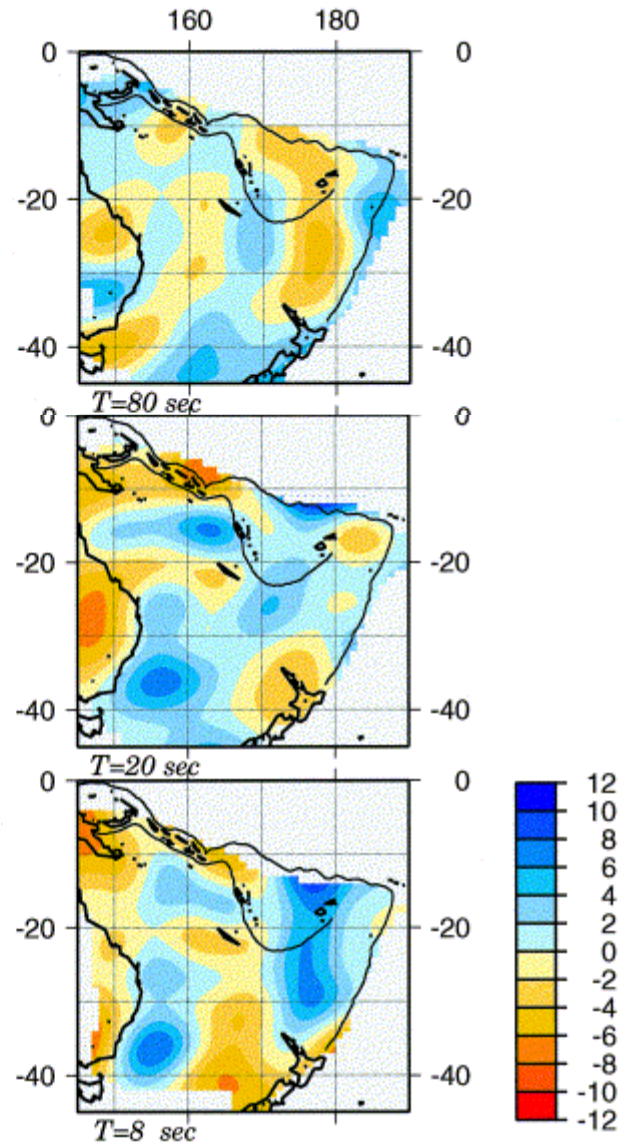


Too much?
Too smooth?

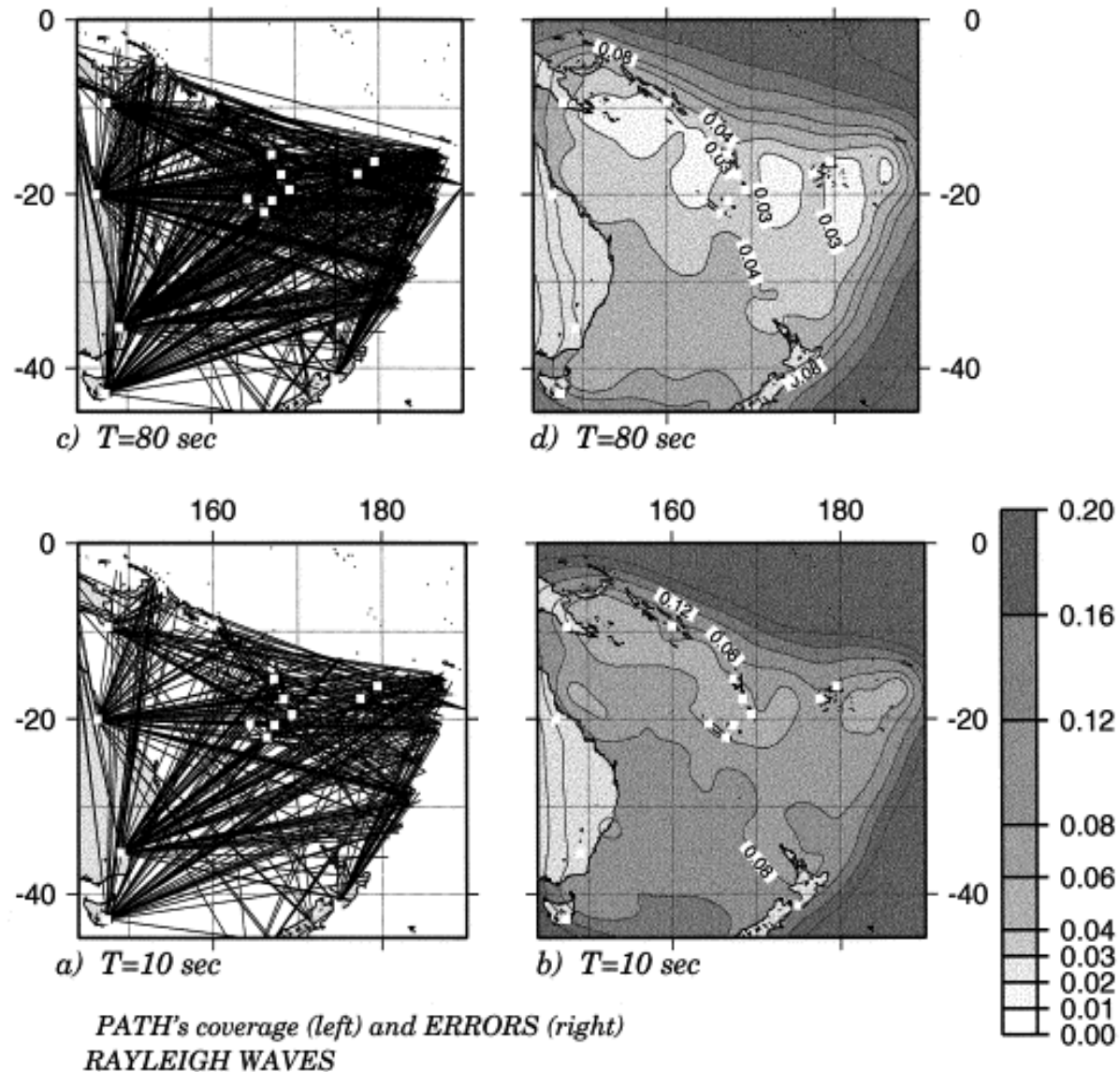


Too little?
Too rough?

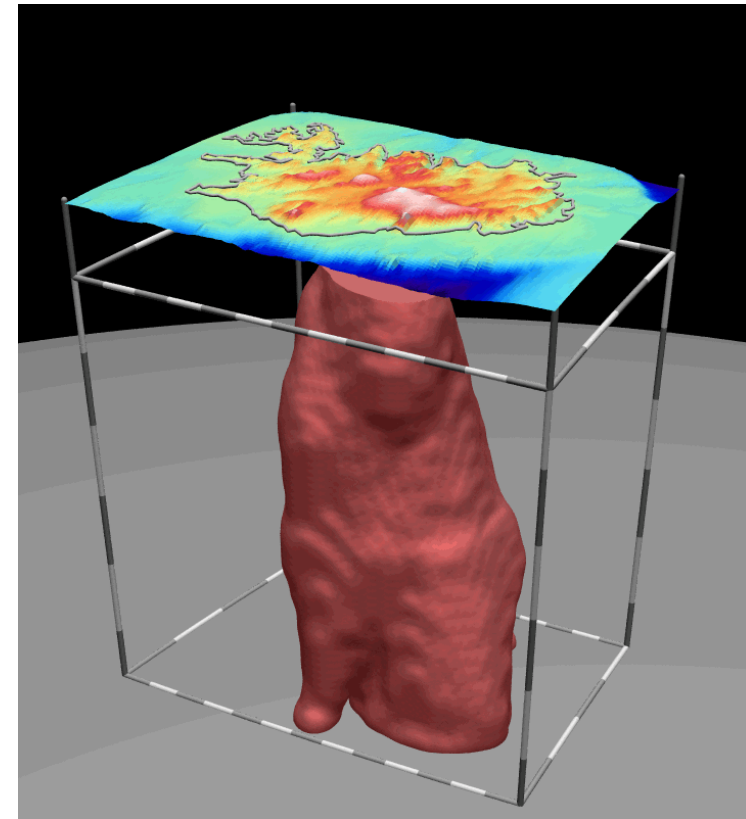
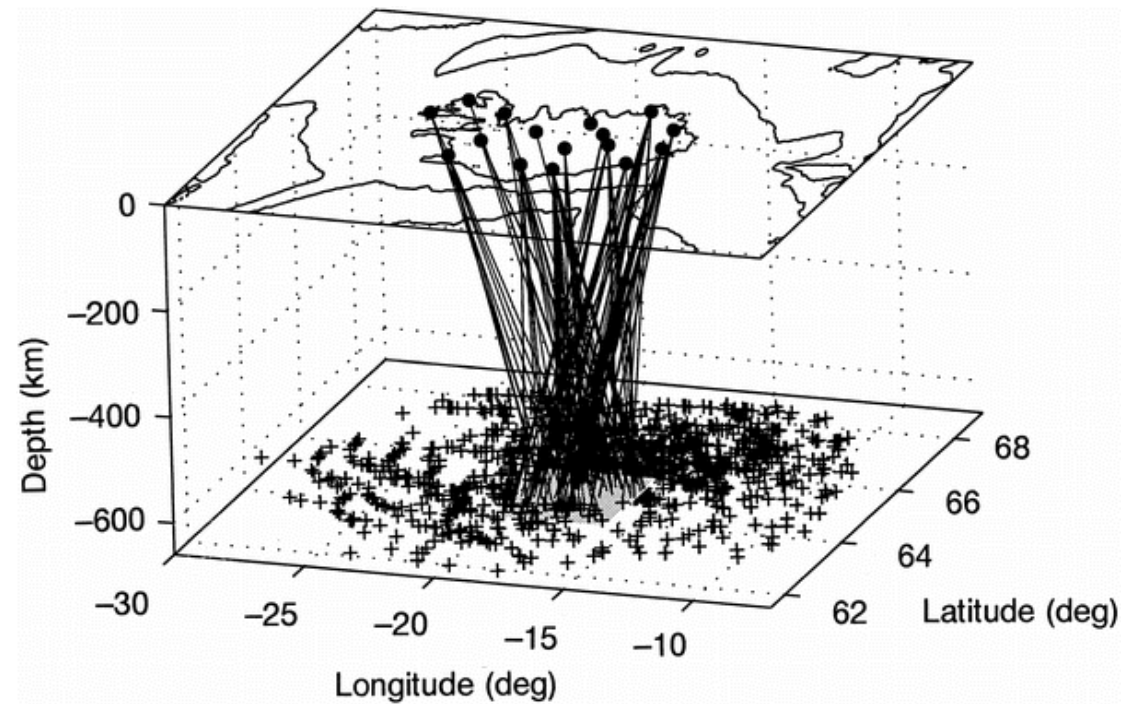
How to interpret seismic models



Demand to see the ray paths



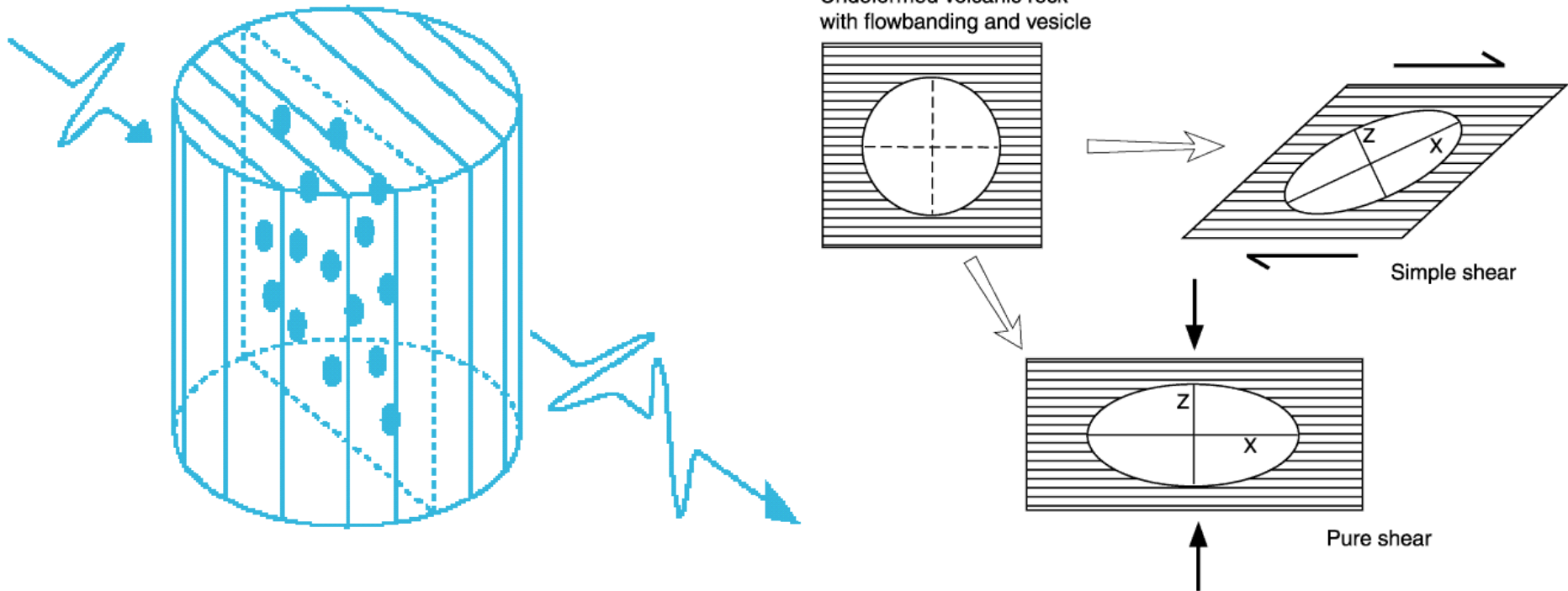
Nature isn't always kind



Seismic anisotropy

Wave speeds depend on

propagation direction and polarization:



No surprise: elasticity maps stress and strain, and both depend on three directions

Polarization anisotropy

- The particles of Love and Raleigh surface waves move in orthogonal directions
- SH and SV body waves sometimes exhibit clear splitting

Azimuthal anisotropy

- It's usually very hard to separate whether the time difference arises from an anisotropic direction or an isotropic wave speed difference (aka heterogeneity)

Why is this so hard?

For 3-D heterogeneity and slight anisotropy:

$$\delta\hat{\beta}_V = \delta\beta_V^{TI} + \frac{G_c}{2\rho\beta_V} \cos 2\theta + \frac{G_s}{2\rho\beta_V} \sin 2\theta \quad (3)$$

Maximum direction is related to fast axis of anisotropic minerals:

$$G = \sqrt{G_c^2 + G_s^2} \quad \text{and} \quad \Psi_{\max} = \frac{1}{2} \arctan \frac{G_s}{G_c} \quad (4)$$

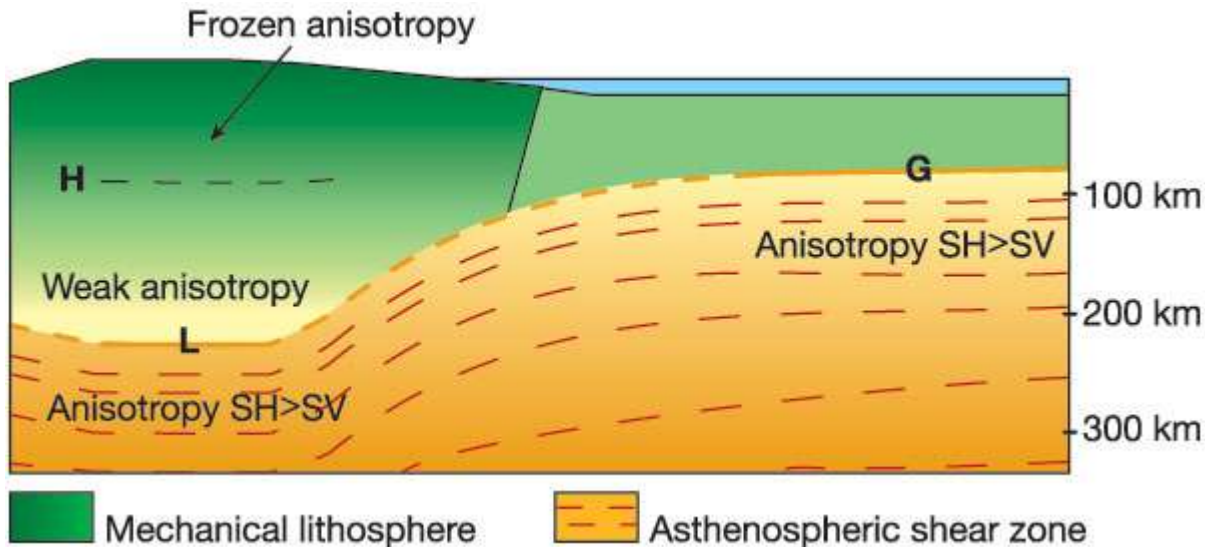
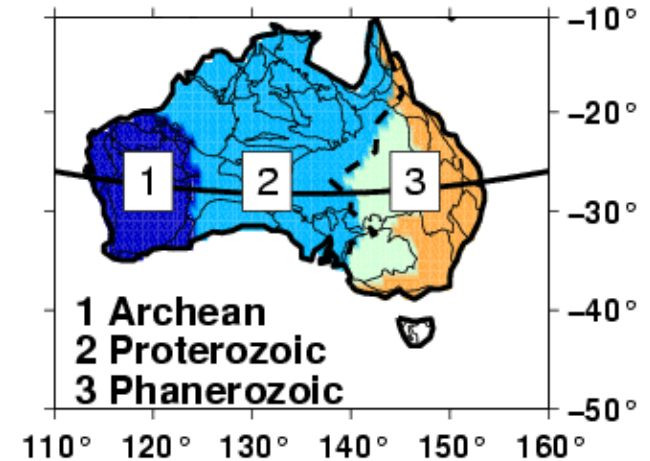
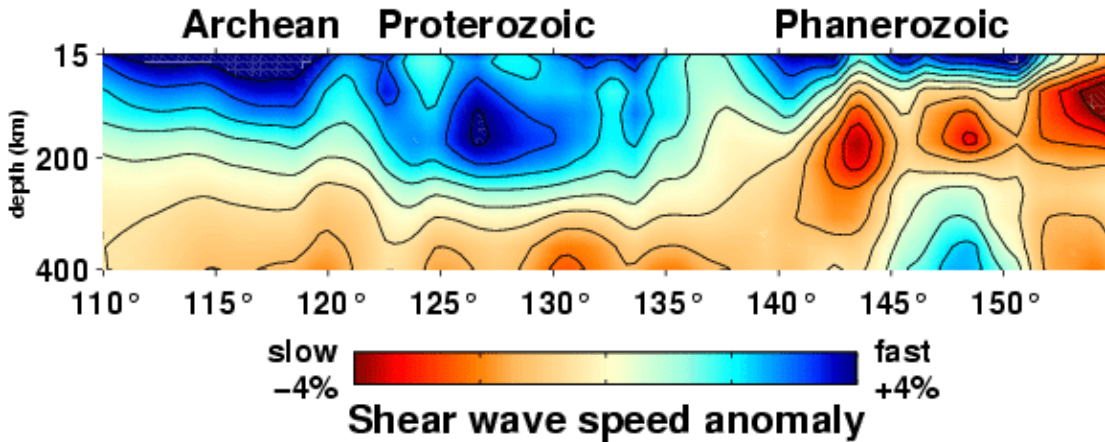
It's very hard to tell whether a phase comes in early because it went through a fast patch or because it came in a fast direction – heterogeneity and anisotropy “trade off.”

Questions to ask of the tomographer

- How is the forward model computed?
- What is the ray coverage?
- What (sort of) damping did you use?
- What does velocity estimation trade off with?
- What is the grid size / the correlation length?
- How are different data sets weighted?
- How far is the final from the starting model?
- Does the starting model have discontinuities?
- How is the surface/depth parameterization
- Is your sensitivity 1-D, 2-D, or 3-D?

Journey to Middle Earth, Part I:

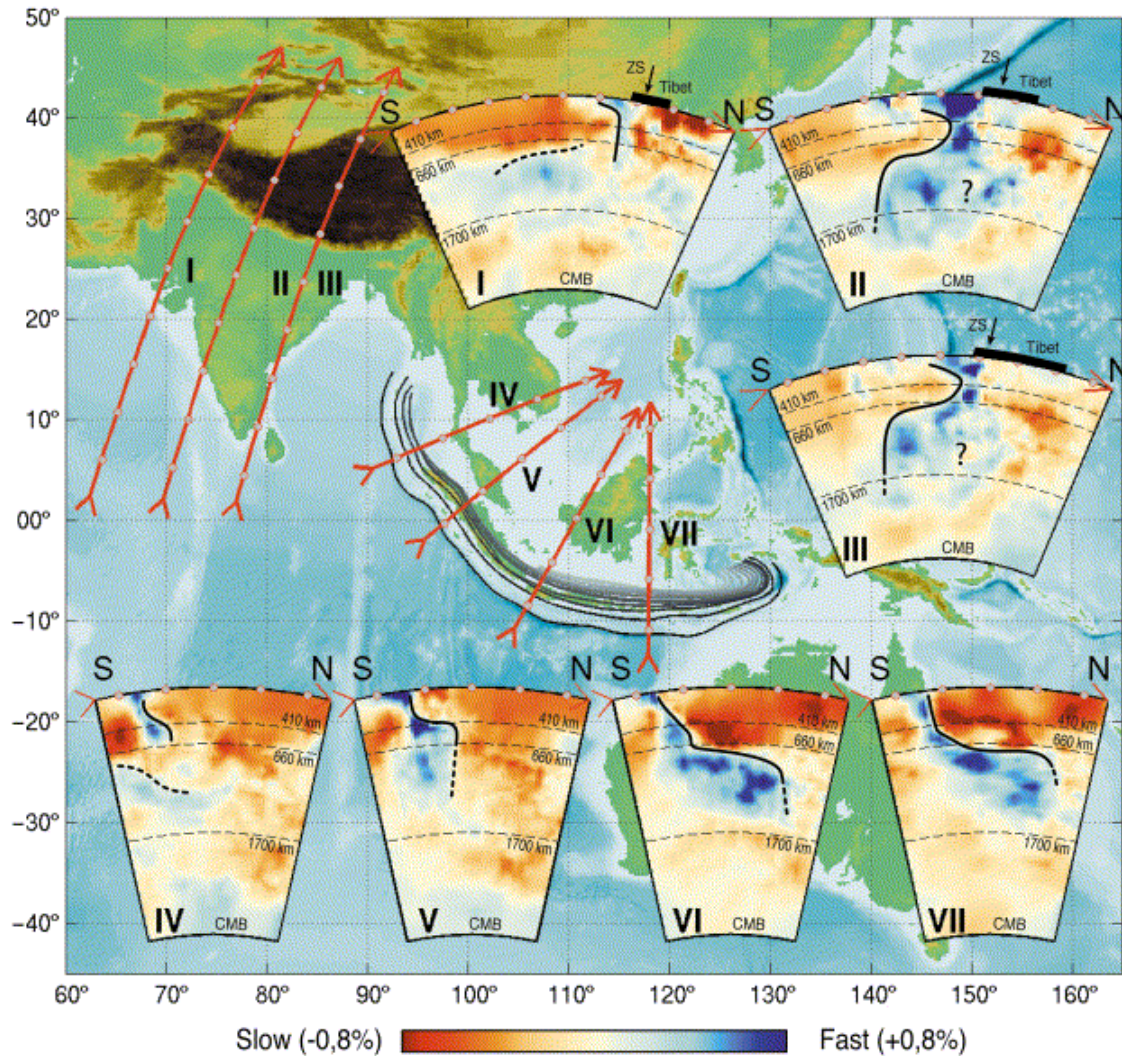
The continental lithosphere



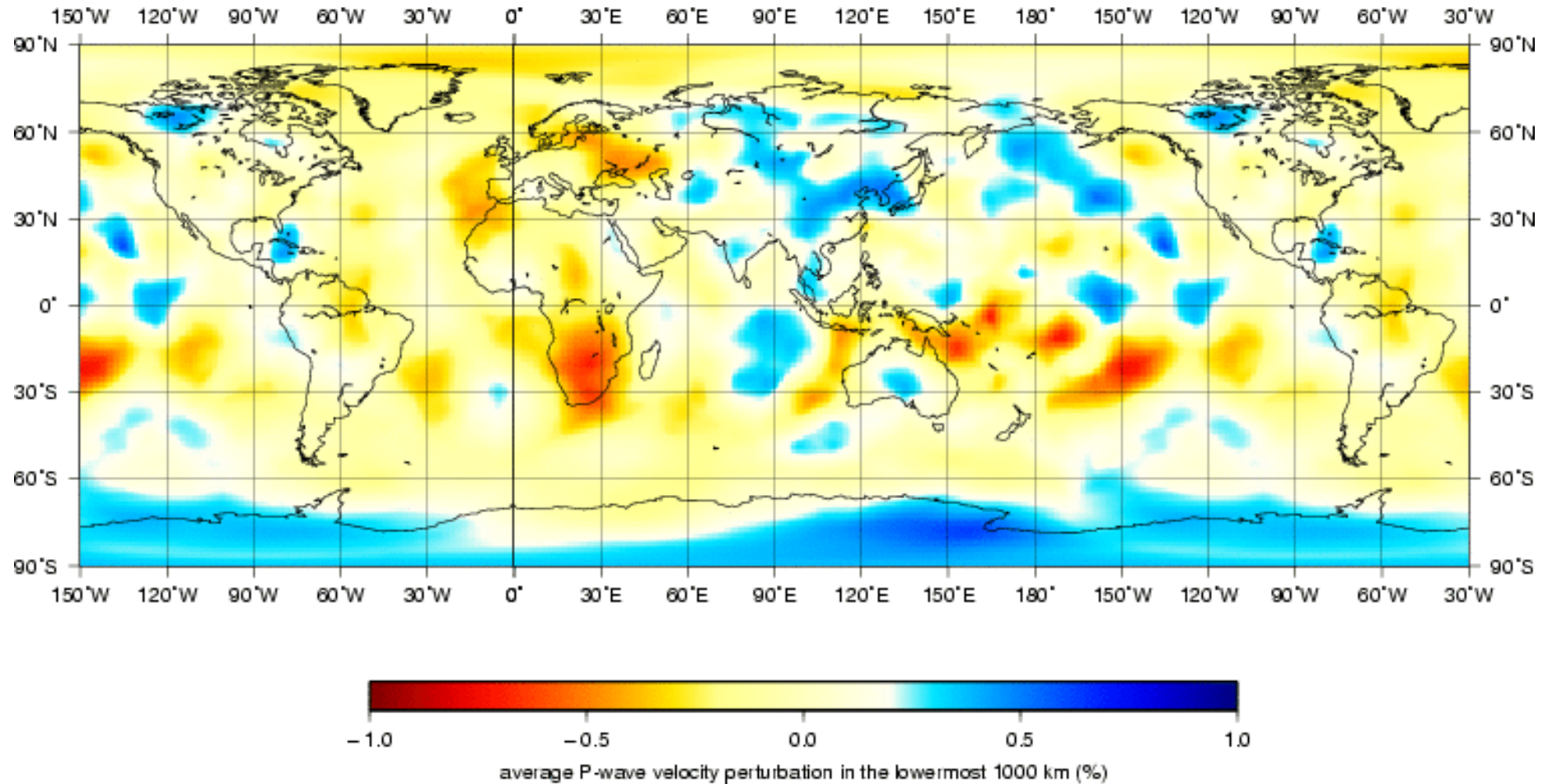
Simons, *GRL*, 2002

Gung, *Nature*, 2003

Journey to Middle Earth, Part II: Subduction zones

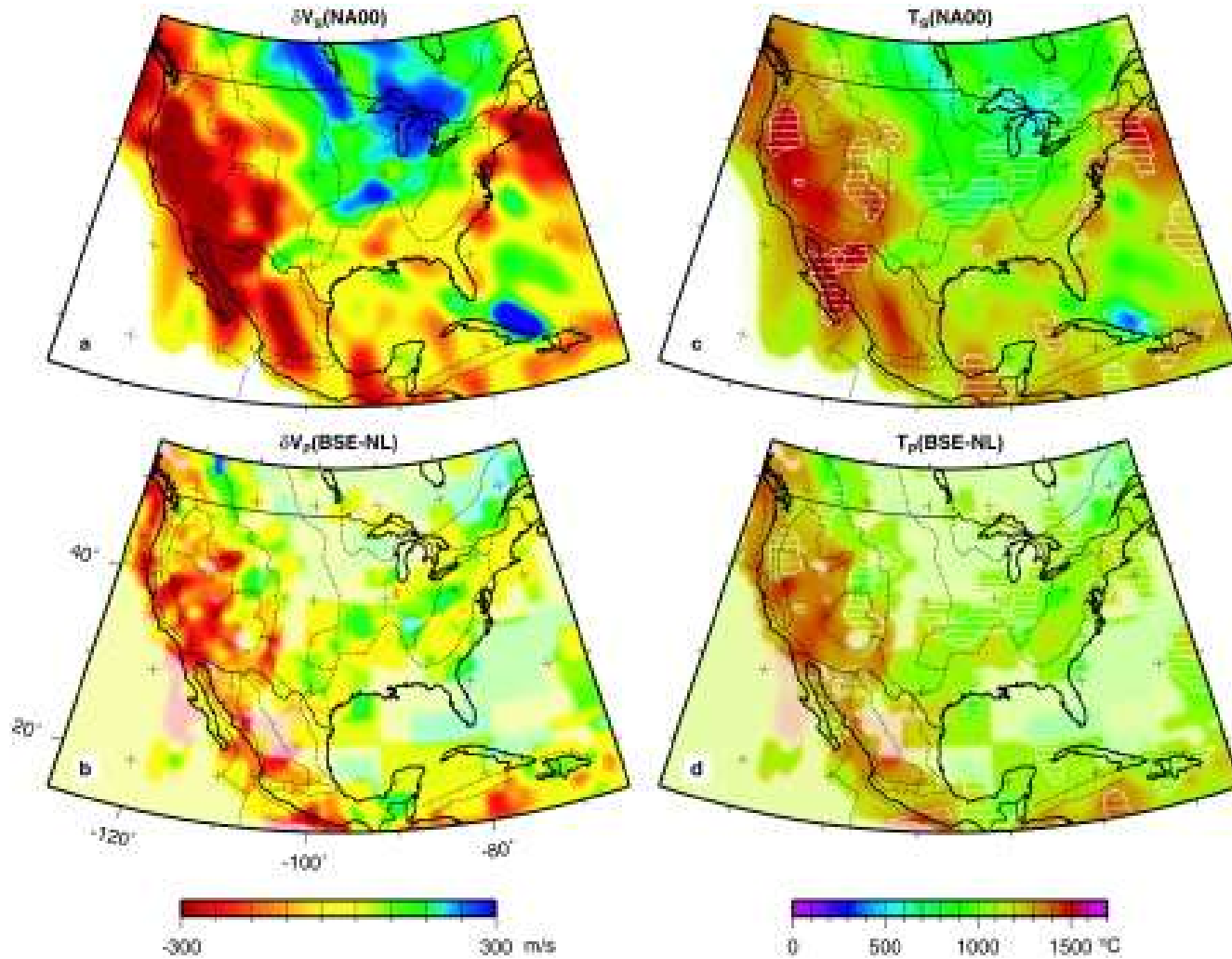


Journey to Middle Earth, Part III: Deep mantle plumes



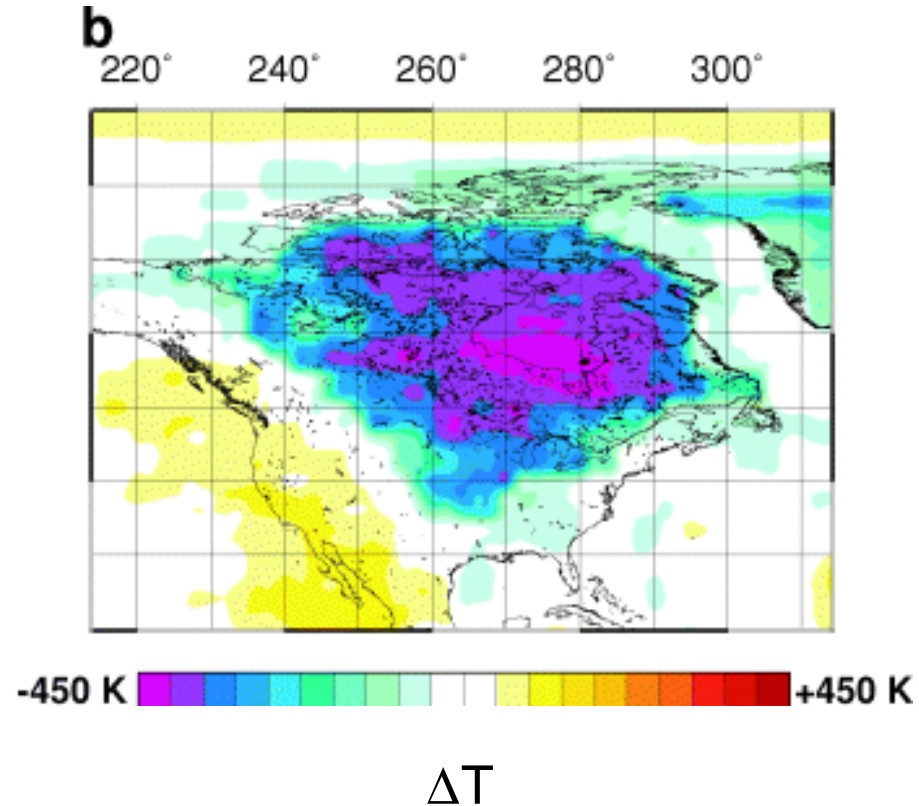
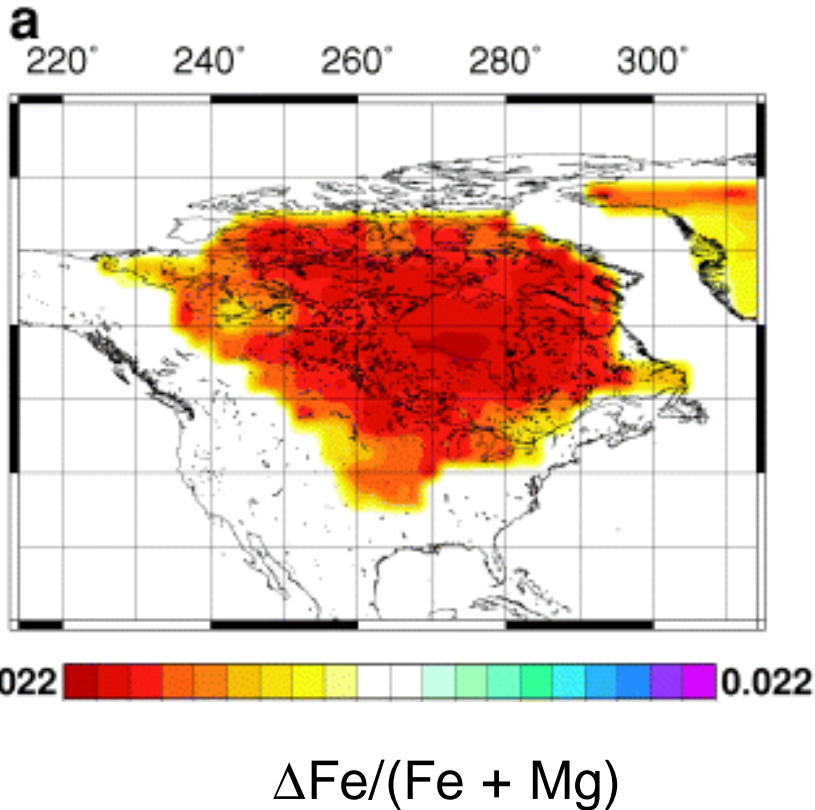
What does it all mean? Part I:

Temperature anomalies



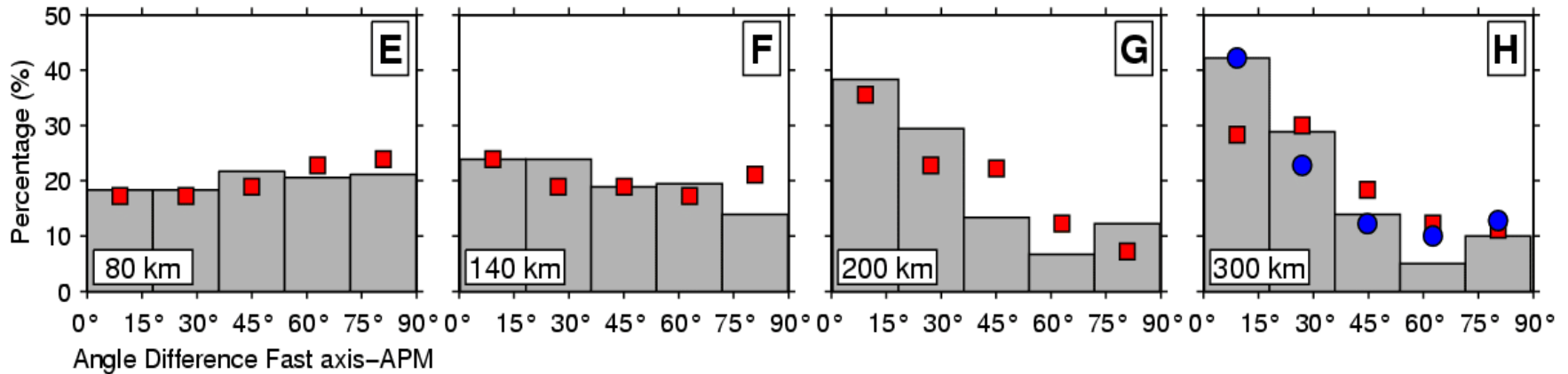
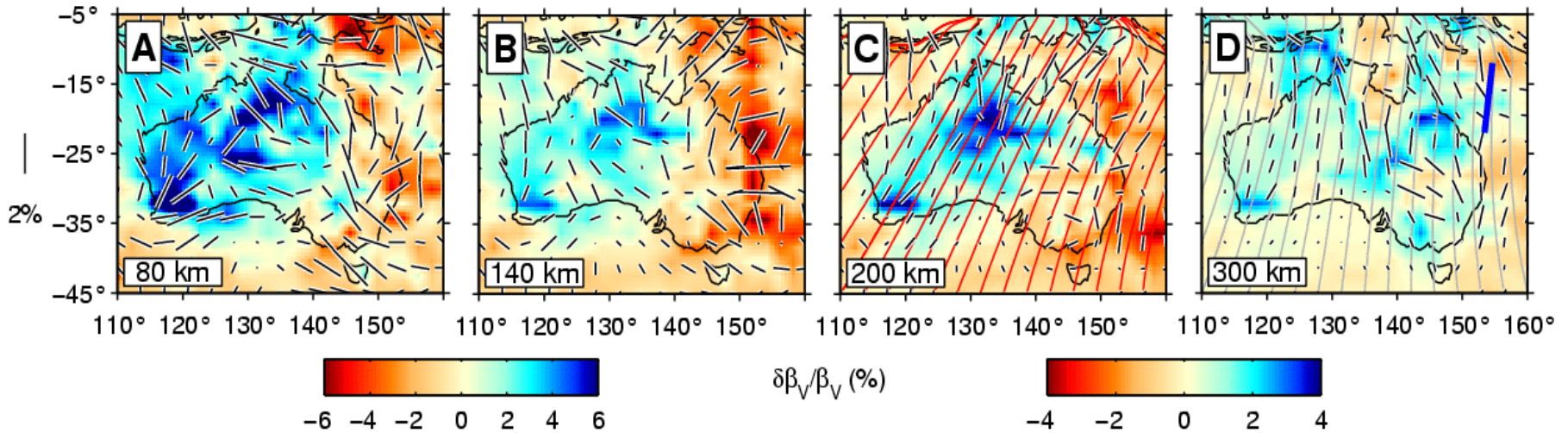
110 km

What does it all mean? Part II: Compositional anomalies



150 km

What does it all mean? Part III: Deformation in the mantle



Fossil

Contemporaneous

Conclusions

- Ultimately, seismology can only tell us where, or in which direction, wave propagation is faster or slower than a reference model
- The non-seismologist has to know the basics of inverse problem modeling, understand the sometimes poor constraints, and be critical
- Improvements are being made: better data, better forward models, better inversions
- As much as with the *a posteriori* interpretation, the community needs to help defining *a priori* acceptable starting models

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**More equations, for
completeness**

A linear system of equations

We're attempting to solve

$$\mathbf{d} = \mathbf{G} \cdot \mathbf{m} \quad (1)$$

Minimize penalty function of weighted error and model norms

$$\begin{aligned} \Phi = & (\mathbf{d} - \mathbf{G} \cdot \mathbf{m}) \cdot \mathbf{A}^{-1} \cdot (\mathbf{d} - \mathbf{G} \cdot \mathbf{m}) \\ & + \mathbf{m} \cdot \mathbf{B}^{-1} \cdot \mathbf{m} \end{aligned} \quad (2)$$

In matrix form, solve

$$\begin{bmatrix} \mathbf{A}^{-1/2} \cdot \mathbf{G} \\ \mathbf{B}^{-1/2} \end{bmatrix} \cdot \mathbf{m} = \begin{bmatrix} \mathbf{A}^{-1/2} \cdot \mathbf{d} \\ 0 \end{bmatrix} \quad (3)$$

Solution

$$\mathbf{m} = (\mathbf{B}^{-1} + \mathbf{G}^T \cdot \mathbf{A}^{-1} \cdot \mathbf{G})^{-1} \cdot \mathbf{G}^T \cdot \mathbf{A}^{-1} \cdot \mathbf{d} \quad (4)$$

Norm and first gradient regularization

For \mathbf{A}^{-1} , use the inverse of the data covariance matrix \mathbf{C}_d (BLUE)

For \mathbf{B}^{-1} , use the identity matrix \mathbf{I} plus the squared first derivative

$$\mathbf{D}_1 = \begin{pmatrix} \dots & 0 & 0 \\ -1 & 1 & 0 \\ 0 & -1 & 1 \\ 0 & 0 & \dots \end{pmatrix} \quad (5)$$

Minimize weighted penalty function

$$\begin{aligned} \Phi = & (\mathbf{d} - \mathbf{G} \cdot \mathbf{m}) \cdot \mathbf{C}_d^{-1} \cdot (\mathbf{d} - \mathbf{G} \cdot \mathbf{m}) \\ & + \alpha \mathbf{m} \cdot \mathbf{I} \cdot \mathbf{m} + \beta \mathbf{m} \cdot \mathbf{D}_1^2 \cdot \mathbf{m} \end{aligned} \quad (6)$$

Solution

$$\mathbf{m} = (\alpha \mathbf{I} + \beta \mathbf{D}_1^2 + \mathbf{G}^T \cdot \mathbf{C}_d^{-1} \cdot \mathbf{G})^{-1} \cdot \mathbf{G}^T \cdot \mathbf{C}_d^{-1} \cdot \mathbf{d} \quad (7)$$

Bayesian inversion

Gaussian *a priori* probability function on the model parameters

$$\rho(\mathbf{m}) \propto \exp\left(-\frac{1}{2}\mathbf{m} \cdot \mathbf{C}_m^{-1} \cdot \mathbf{m}\right) \quad (8)$$

Maximize joint distribution of data, model, subject to $\mathbf{d} = \mathbf{G} \cdot \mathbf{m}$

$$\boxed{\mathbf{m} = \mathbf{C}_m \cdot \mathbf{G}^T \cdot (\mathbf{C}_d + \mathbf{G} \cdot \mathbf{C}_m \cdot \mathbf{G}^T)^{-1} \cdot \mathbf{d}} \quad (9)$$

Equivalent to (using a trivial matrix identity)

$$\mathbf{m} = (\mathbf{C}_m^{-1} + \mathbf{G}^T \cdot \mathbf{C}_d^{-1} \cdot \mathbf{G})^{-1} \cdot \mathbf{G}^T \cdot \mathbf{C}_d^{-1} \cdot \mathbf{d} \quad (10)$$

So in choosing norm and gradient regularization we've identified

$$\mathbf{C}_m^{-1} = \alpha \mathbf{I} + \beta \mathbf{D}_1^2 \quad (11)$$

This imposes a particular form of the *a priori* covariance \mathbf{C}_m

To Bayes or not to Bayes, what's the question?

A priori model covariance function with correlation length L

$$C_{\mathbf{m}}(\mathbf{r}_1, \mathbf{r}_2) = \sigma^2 \exp\left(-\frac{|\mathbf{r}_1, \mathbf{r}_2|^2}{2L^2}\right) \quad (12)$$

The following equivalence holds [*Yanovskaya and Ditmar, 1990*]

$$\mathbf{m} \cdot \mathbf{C}_{\mathbf{m}}^{-1} \cdot \mathbf{m} = \frac{1}{2\pi} \frac{1}{(\sigma L)^2} \sum_{n=0}^{\infty} \frac{1}{n!} \left(\frac{L^2}{2}\right)^n \nabla^n \mathbf{m} \cdot \nabla^n \mathbf{m} \quad (13)$$

So indeed

$$\mathbf{C}_{\mathbf{m}}^{-1} = \alpha \mathbf{I} + \beta \mathbf{D}_1^2 + \text{higher-order terms} \quad (14)$$

Exact resolution computation

For the linear problem, in a generalized sense,

$$\mathbf{m}^{\text{est}} = \mathbf{G}^{-\text{g}} \cdot \mathbf{d}^{\text{obs}} = \mathbf{G}^{-\text{g}} \cdot \mathbf{G} \cdot \mathbf{m}^{\text{true}} \quad (15)$$

The resolution matrix is given by

$$\mathbf{R} = \mathbf{G}^{-\text{g}} \cdot \mathbf{G} \quad (16)$$

In the Bayesian framework [*Montagner, 1986*]

$$R(\mathbf{r}, \mathbf{r}') = \delta(\mathbf{r} - \mathbf{r}') - \frac{C_{\text{p}}(\mathbf{r}, \mathbf{r}')}{C_{\text{m}}(\mathbf{r}, \mathbf{r}')} \quad (17)$$

This represents the degree to which we are able to reduce the *a priori* covariance C_{m} of the model parameters (the null-state of information) by obtaining the *a posteriori* covariance structure C_{p} after the inversion.