Geodynamics III: Integrating geophysical observations into global mantle flow models

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Composite by D. Evans & Co., showing geological unit compilation by D. Eglington, seafloor age from Mueller et al. (2008), SAVANI tomography by Auer et al. (2014)

Dziewonski plot



reconstructed seafloor age

60'1

60'

30"

0 Ma

100 Ma

hot temperature isosurface

dense chemical piles





Olson et al. (2016)











Eurasi

Seton et al. (2012)

cf. McNamara and Zhong (2005) and many others



We know how plate tectonics works (~true, and often assumed)

Plate tectonics is the top boundary layer of thermo-chemical convection



Symmetry of upwellings and downwellings of Rayleigh-Taylor convection is broken by:

- Image: temperature dependence of viscosity
- > depth-dependent viscosity
- → internal vs. bottom heating
- → fractionation (e.g. continents and thermo-chemical piles)

→ ...

Plate driving forces: Integrals over individual components of thermo-chemical convection



Lithospheric thickening (AKA ridge push, oceanic GPE)





Slab pull (subducting thermal boundary layer)



Force estimates from half-space cooling (hugely important reference model and

achievement of geodynamics)

- Ridge push (lithospheric thickening) $\sim 10^{12}$ N/m
- Slab pull ~ 10^{13} N/m





FIG. 8. Percentage of circumference of plate connected to downgoing slab. Open bar is total length, filled bar is effective length.

Forsyth & Uyeda (1975)

How does the mantle drive the plates at present? Estimate from global circulation modeling



Forte (2015)



Yield stress (MPa)

cf. Ricard and Vigny (1989), Forte (1993), King and Hager (1990), Gable et al. (1991), Han and Gurnis (1999) Alisic et al. (2012)

Why do we have plate tectonics?



free-slip, global convection computation with temperature-dependent viscosity and yield stress, $Ra > Ra_{Earth}$

van Heck and Tackley (2008); Yoshida (2008); Foley and Becker (2009); Nakagawa et al. (2009); Coltice et al. (2013), Coltice et al. (2016)



 $\delta v(\theta, \phi) \approx \sum_{l=0}^{\ell_{\text{max}}} \left[a_{\ell 0} X_{\ell 0}(\theta) + \sqrt{2} \sum_{m=1}^{\ell} X_{\ell m}(\theta) \times (a_{\ell m} \cos m\phi + b_{\ell m} \sin m\phi) \right]$



Stochastic constraints: Mantle tomography shows long wavelength (*I* ~ 2) structure

mantle power spectrum from seismology (S20RTS) 500 1000 depth [km] 1500 2000 2500 15 10 20 5 degree I normalized power 0.25 0.50 0.75 1.00

Convection models predict this heterogeneity

(because it has plates and those organize convection)



cf. Tackley (2000a,2000), Richards et al. (2001), van Heck and Tackley (2008), Mallard et al. (2016)

We don't know how plate tectonics works

(questioning models assuming known physics)

Seafloor age distribution vs. boundary layer instability

age α from Muller et al. (2008) 3 triangular fit age α [km²/yr] area per 0 50 100 150 seafloor age τ [Myr]

cf. Becker et al. (2009) Coltice et al. (2013) 200

Kinematic vs. dynamic models



Bower et al. (2014)



Geophysical constraints on global mantle dynamics (applied geodynamics)

- Plate velocities
- Topography
- Geoid
- Seismic tomography
- Seismic anisotropy



Argus et al. (2011), Becker et al. (2015)

Helmholtz decomposition

into irrotational (poloidal) and divergence free (toroidal) fields

$$\underline{v}_{h} \equiv \underline{v}_{h,\text{pol}} + \underline{v}_{h,\text{tor}} \equiv \underline{\nabla}_{h} \varphi + \underline{\nabla} \times (\psi \hat{z})$$
$$\nabla^{2} \varphi = \underline{\nabla}_{h} \bullet \underline{v}_{h}; \nabla^{2} \psi = \text{vort}_{z}(\underline{v}_{h})$$

foroidal



Or, in *vector* spherical harmonics (use *generalized* spherical harmonics for orientational, 2φ):

$$\mathbf{V}(\theta,\phi) = \sum_{l\ m} Re\{b_{lm}\mathbf{B}_{lm}(\theta,\phi) + c_{lm}\mathbf{C}_{lm}(\theta,\phi)\}$$
$$\mathbf{B}_{lm}(\theta,\phi) = \frac{1}{L}\frac{\partial P_{lm}(\theta)}{\partial \theta}e^{im\phi}\mathbf{e}_{\theta} + \frac{im}{L\sin\theta}P_{lm}(\theta)e^{im\phi}\mathbf{e}_{\phi}$$

$$\mathbf{C}_{lm}(\theta,\phi) = \frac{im}{L\sin\theta} P_{lm}(\theta) e^{im\phi} \mathbf{e}_{\theta} - \frac{1}{L} \frac{\partial P_{lm}(\theta)}{\partial \theta} e^{im\phi} \mathbf{e}_{\phi}$$



time [Ma]

Surface topography



Isostasy



p = const. @ compensation level



crustal thickness



Note: very uneven coverage. active source best, RF next best thing

actual topography



crustal thickness from CRUST1.0



$$f_{1} = \frac{\rho_{a} - \rho_{c}}{\rho_{a}} \quad f_{2} = \frac{\rho_{a} - \rho_{l}}{\rho_{a}}$$

$$f_{1} = 0.12$$

$$f_{2} = -0.01$$

$$f_{2} = -0.01$$

$$f_{2} = -0.01$$

$$f_{3} = -0.01$$

$$f_{4} = -0.01$$

$$f_{2} = -0.01$$

$$f_{3} = -0.01$$

£1 | £1

2

observed topography [km]

10 20 30 40 50 60 70 effective crustal thickness [km]

0.002

0.001

actual topography



crustal thickness from CRUST1.0



$$\hat{t} = f_1 l_c + f_2 l_l$$

$$f_1 = \frac{\rho_a - \rho_c}{\rho_a}$$
 $f_2 = \frac{\rho_a - \rho_l}{\rho_a}$



topography



Global crustal residual topography

crustal model



Note: Like oh so many things, "model", not "data"

Non-isostatic, ~Airy residual topography



Geoid anomalies



(corrected for *hydrostatic* shape)



Potential field constraints





Figure from Y. Ricard

How do we build a deep Earth density model?



Upper mantle slabs from Wadati-Benioff zones



Hager (1984)


Stokeslets link subduction history to mantle structure



Ricard et al. (1993); Lithgow-Bertelloni et al. (1993, 1998)

Seismic tomography shows slabs in lower mantle

smean @ 1200 km



SLABS (DEPTH 2000 KM, DEGREES 1-15)



SLABS (DEPTH 2000 KM, DEGREES 1-3)



SH425.2 (DEPTH 2000 KM, DEGREES 1-3)



CORRELATION SLABS/SH425.2



Ricard et al. (1993)

Mantle structure as *f*(*t*) from subduction





Ricard et al. (1993); Lithgow-Bertelloni & Richards (1998); Steinberger (2000); Spasojevich et al. (2009); Steinberger and Torsvik (2010) Steinberger et al. (2014); Bower et al. (2015)

Possible reference frame (but sinking rates uncertain)



van der Meer et al. (2009)

Correlation: Advected slabs vs. tomography

stb00d vs. smean, $\langle r_{20} \rangle = 0.21$, $\langle r_8 \rangle = 0.30$



Lithgow-Bertelloni & Richards (1998); Steinberger (2000); Bunge & Grand (2000); Spasojevic & Gurnis (2009)

Becker & Boschi (2002)



Note: Correlation with moving plumes better than with slabs

Boschi et al. (2007, 2008)

Advected plumes vs. tomography (Steinberger flow model conduits)



Boschi et al. (2007, 2008)

Building a model of Earth from slabs and plumes

Input: Slabs and plumes

z = 500 km. input



z = 1000 km, input







z = 2500 km, input



Resolving power of tomography

z [km]



Nataf & Ricard (1993); Megnin et al. (1997); Boschi (2003); Ritsema et al. (2007)

Inter tomography-model correlation



Note: data and theory matter

Auer (2016)



How to convert tomography to density anomalies? I. Take it from mineral physics and composition from ..

Stixrude and Jeanloz (2009)



II. Best-fit scaling and composition as inverse problems



cf. Simmons et al. (2008, 2010), Soldati et al. (2014), Forte et al. (2015)

Besides cratons and piles: dln ρ /d ln $v_s = 0.2$

∆T @ 1200 km



 Use mineral physics to convert velocity into temperature (density) anomalies

How to compute flow?

Mantle circulation

- Treat mantle and lithosphere as a fluid
- Infinite Prandtl number (no inertia) approximation
 - Navier-Stokes turns into Stokes equation
- Instantaneous solution for given density and boundary conditions
 - Can solve in <~1 s for spherical Earth without lateral viscosity variations (LVVs)





Asides:

- \rightarrow A needle, as opposed to sphere, will sink with 0.5...2 V_{Stokes}
- → Pe # = ratio of diffusive to convective time scale = t_{d} / t_{c}

 $t_c = a^2 / \kappa$, $t_c = a / v_{\text{Stokes}}$, $\Delta \rho = \Delta T \alpha \rho_0$, then Pe \rightarrow Ra (with *a* instead of *L*)

Mantle circulation

- Thermo-chemical heterogeneity and complex rheologies make things interesting
- Finite element methods best suited for lateral viscosity variations (we can now solve all of this, at < km resolution without approximations)



force balance (conservation of momentum)

> constitutive law (rheology)

$$\frac{\partial \sigma_{ij}}{\partial x_{j}} = -Ra \Delta T \delta_{ir} - Ra_{C} \Delta C \delta_{ir}$$

$$chemical buoyancy$$

$$\sigma_{ij} = -P \delta_{ij} + 2\eta (\sigma, T, d, H_{2}O, \varepsilon) \delta_{ij}$$

non-Newtonian viscosity with memory

Mantle convection

- * Energy equation introduces time-dependence
- Coupling between velocity and temperature introduces non-linearity
 - * Can time reverse advection, but not diffusion



Global circulation models

show that plate-motion induced shear cannot always be guessed from surface motions



Let's fit the geoid: Static effect of slablet



sketch from M. Billen's MYRES talk cf. Hager & Richards (1984); Ricard (1984)

Combined static and dynamic geoid effect of slablet



sketch from M. Billen's MYRES talk cf. Hager & Richards (1984); Ricard (1984)

Kernels for homogeneous mantle





Colli et al. (2016), sketch from M. Billen cf. Hager & Richards (1984); Ricard (1984)

Geoid for tomography driven flow

Isoviscous – free slip surface boundary condition



Geoid for tomography driven flow

 Four layer model with viscosity increase in lower mantle



"Airy" residual surface topography: mainly due to half-space cooling



Mueller et al. (2008) seafloor ages

0

→ the outstanding performance of this geodynamic "model" complicates attribution of anomalous topography

Non half-space cooling residual topography





 \rightarrow dynamic topography *h* μ with density (temperature anomaly), $\Delta \rho$ (here = ΔT)

Stokes' sphere velocity/stress:

 $v_{Stokes} \mu \Delta \rho \eta$ $\tau \mu \Delta \rho$ Notes:

- \rightarrow Often, we infer equivalent topography from
 - $\sigma_{_{\! ZZ}}$ pushing on a free slip surface

this works remarkably well in most cases ($\lambda \ll L$)

Dynamic topography physics: "plume" case



 \rightarrow dynamic topography *h* μ with density (temperature anomaly), $\Delta\rho$ (here = ΔT)

 $v \mu \Delta \rho / \eta$ h μ τ μ Δρ dh/dt μ v h μ Δρ²/η

Note:

- Match topography and velocity, constrain both Δρ and η,
- match uplift, even better!



 \rightarrow uplift rate, dh/dt μ inverse of viscosity, 1/ η (and density anomaly squared)



Isochemical, $\Delta \rho = 1$



0.0 0.2 0.4 0.6 0.8 1.0 Temperature/Composition

dynamic topography @ CMB



Lassak et al. (2008)

Pilotania, $\Delta \rho = 1$



Lassak et al. (2008)

Pilotania, $\Delta \rho = 1,000$



Lassak et al. (2008)

Viscosity inversions are non-unique (Monte Carlo approach, based on geoid and surface dyn. topo)



Panasyuk & Hager (2000)

Combining post-glacial rebound (GIA) and geoid (still need to close loop with ice models)

Notes:

- geoid gives
 relative η with depth
- →GIA gives absolute value
- → "Haskell constraint" is ~10²¹ Pas for average down to ~1200 km (under cratons...)


Slab ponding diversity and viscosity stratification in transition zone



Nondimensional viscosity

Global, stochastic view: Radial correlation functions



cf. Puster and Jordan (1997), Tackley (1998, 2002)

Query data for discontinuity depth

relative variance reduction



cf. Puster and Jordan (1997), Tackley (1998, 2002)

see Max Rudolph talk on Rudolph et al. (2015)

Adding a constraint: Predicting plate motions

Flow model with only radial viscosity variations Poloidal component



- No toroidal flow without lateral viscosity variations (no PT coupling)
- Strain-rates not very plate-like

Observed plate velocities in hot spot reference frame Poloidal component Toroidal component



O'Connell et al. (1991); Ricard et al. (1991); Ribe (1993); Forte & Peltier (1993); Thoraval & Richards (1997); Moucha et al. (2008); Ghosh et al. (2009)

Rigid plate motions for weak boundaries



 $\mathbf{T}_{i} = \int_{\text{plate}} dA \ \mathbf{r} \times \sigma$ $\mathbf{T}_{vd} = \mathbf{P} \cdot \boldsymbol{\omega} = \sum_{i}^{M} \mathbf{T}_{i}$

Note: Can construct plate motion interaction matrix **P**, for any LVVs

cf. Ricard and Vigny (1989), Forte (1993), van Summeren et al. (2010)

Match to observed plate motions

- Velocity model
 - Prescribe weak plate boundaries
 - Compute plate drag coupling and driving torques
 - Solve for Euler vectors for rigid plates
- Correlations good, but oceanic plates move as fast as continental ones

Observed plate motions



Ricard & Vigny (1989); Lithgow-Bertelloni & Richards (1998); Becker & O'Connell (2001); Conrad & Lithgow-Bertelloni (2002); Becker (2006); van Summeren et al. (2010; Alisic et al. (2014) Ricard & Vigny (1989)



Becker & O'Connell (2001)

Ricard & Vigny (1989); Lithgow-Bertelloni & Richards (1998); Zhong et al. (2000); Becker & O'Connell (2001); Conrad & Lithgow-Bertelloni (2002); Becker (2006)



Conrad and Lithgow-Bertelloni (2004)

Mantle rheology

- use generic temperature and/or stress dependence of viscosity on top of radial variations
- use effective, olivine creep law with diffusion & dislocation creep



 η : viscosity

ε: strain-rate

$$\eta = \left(\frac{d^m}{A'}\right)^n \dot{\varepsilon}_{II}^{\frac{1-n}{n}} \exp\left(\frac{E^* + pV^*}{nRT_r(T_c + T)}\right)$$

Sub-oceanic vs. continental speeds (LVV GCM @ 250 km)



Hager & O'Connell (1981); Ricard & Vigny (1989); Zhang & Christensen (1993); Cadek & Fleitout (2003); Becker (2006);

Stadler et al. (2010); Alisic et al. (2014)



van Summeren et al. (2012)

Average mantle viscosities for olivine



Cadek & Fleitout (2003); McNamara et al. (2003);

Podolevsky et al. (2005); Becker (2006)

 $\log_{10}(\epsilon_{disl}/\epsilon_{diff})$

- Dry olivine (Hirth & Kohlstedt, 2003), grain size ~5 mm
- viscosity broadly consistent with geoid and post-glacial rebound
- Depth range of dislocation creep consistent with anisotropy observations

region of

PO formation





b) SL2013SVA vs. NNR @ 200 km





d) SL2013SVA vs. LPO @ 200 km



Becker et al. (2014), using SL2013SV





b) DR2015 vs. NNR @ 200 km







using Debayle et al.'s (2015) model







slab South American craton 10°

retreating

Antilles

5°

Growdon et al. (2009); Masy et al. (2011)

SKS splitting from reference model (tomography driven)



Miller and Becker (2012)

Splitting from slab model with a weak asthenosphere and keel



Miller and Becker (2012)

Inverse geodynamics: Regional splitting azimuth misfit



How to resolve some of the non-uniqueness and uncertainties?





Geodynamic inversions



Baumann and Kaus (2015)



State of affairs

- Can explain plate velocities, geoid, dynamic topography and ~seismic anisotropy with global mantle circulation models
- Provides constraints on rheology and effective density distribution, needed to understand terrestrial planet evolution
- Frontiers: Time evolution and continental dynamics, e.g.
 - predict intraplate deformation
 - experimental design/hypothesis testing

Additional slides



Long and Becker (2010)





Residual topography models

crust for sparse, oceanic sites from active source, corrected to plate model, *continental areas from free-air*

> crust from CRUST1, lithospheric Model, corrected for half-space cooling

Admittance != constant



(CIDER supported...)

Colli et al. (2016) cf. Hager & Richardson (1984); Ricard (1984)

gravity dynamic



Prediction of plate scores with LVVs







Geoid for tomography driven flow

• Lower mantle stiffer





Morgan (1968)




Burov et al.

Alternative age distributions for constant production rates

Subduction probability

Age distribution



Yield stress control



cf. Tackley (2000a,2000), Richards et al. (2001), van Heck and Tackley (2008), Foley and Becker (2009)

But: effect of asthenosphere, internal vs. bottom heating, Ra #, continents, and damage/memory



Two density sources

Free-slip solution for poloidal flow

"Plate" solution for poloidal flow

Forte (2008, 2015)



plate motions plate motions at cratons plate motion induced drag transform fault shear colliding resistance subduction resistance upper mantle slab suction slab pull gravitational sliding ridge push mantle drag slab suction transform fault normal

Geoid and LVVs

- Free slip surface
- Four layer viscosity (not optimized)
- SMEAN tomography







Geoid and LVVs

- Free slip surface
- Weak zones, stiff keels



Surface velocities with viscosity in background









Subduction velocity scaling slab mantle bending drag

slab $V' \approx \Delta \rho g h L / [2\eta_l (h/r)^3 + 3\eta_m A]$ pull $V^\circ = (\Delta \rho^\circ g H^\circ h^\circ) / \eta_m^\circ$ $V = V' / V^\circ$

 Subduction velocity = modified Stokes velocity accounting for bending

Regional models



Becker & Faccenna (2009)

Global models





Gravitational potential energy



Figure 3. Differential potential energy $\Delta U = U - \overline{U} (\overline{U} = 2.616 \times 10^{14} \text{ Jm}^{-2})$ of our 2° × 2° model in units of 10¹⁴ Jm⁻². Min/mean/max values of ΔU for oceanic and continental lithosphere are -0.03/-0.004/ 0.023 and -0.013/0.007/0.078, respectively.

Becker and O'Connell (2001)

GPE torques



Becker and O'Connell (2001)

